

Compatibility Conditions for the Left Cauchy Green Tensor Field in 3-D

Amit Acharya

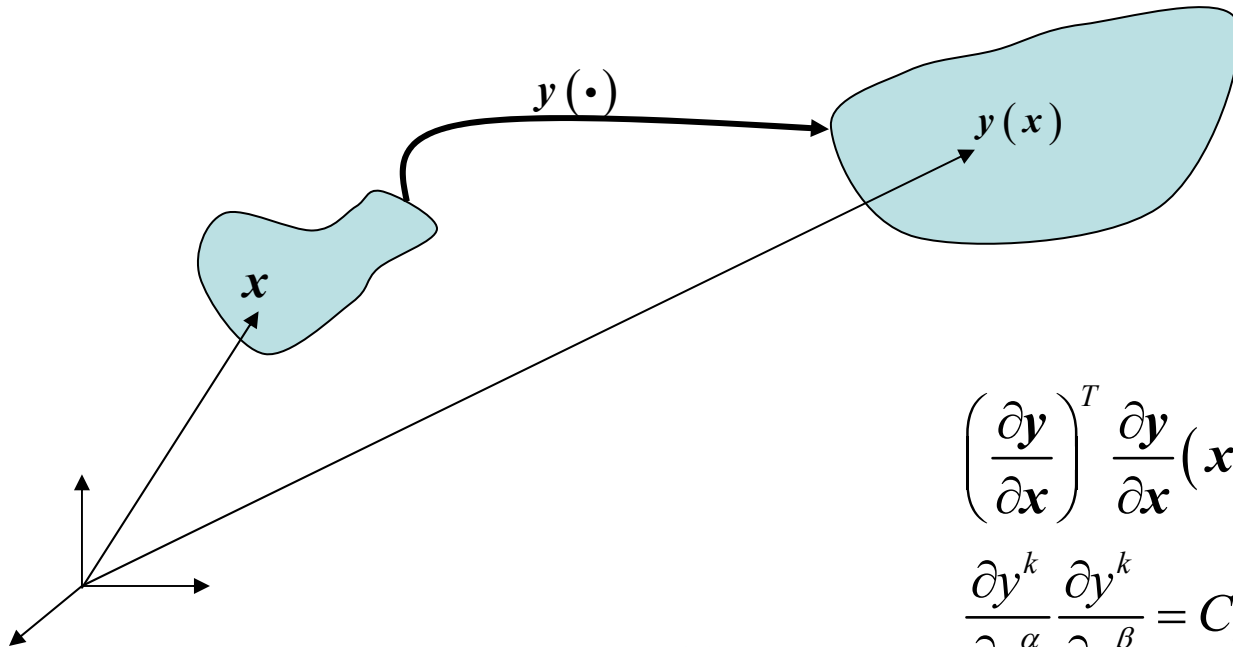
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Continuum Mechanics Problem for Left/Right Cauchy Green Compatibility (LCG/RCG)



Prescribed $\mathbf{C}(\cdot), \mathbf{B}(\cdot)$

as functions of \mathbf{x}

Find $\mathbf{y}(\cdot)$

$$\left(\frac{\partial \mathbf{y}}{\partial \mathbf{x}} \right)^T \frac{\partial \mathbf{y}}{\partial \mathbf{x}}(\mathbf{x}) = \mathbf{C}(\mathbf{x}) \in P_{sym} \quad \text{RCG}$$

$$\frac{\partial y^k}{\partial x^\alpha} \frac{\partial y^k}{\partial x^\beta} = C_{\alpha\beta}$$

$$\frac{\partial \mathbf{y}}{\partial \mathbf{x}} \left(\frac{\partial \mathbf{y}}{\partial \mathbf{x}} \right)^T(\mathbf{x}) = \mathbf{B}(\mathbf{x}) \in P_{sym} \quad \text{LCG}$$

$$\frac{\partial y^i}{\partial x^\alpha} \frac{\partial y^j}{\partial x^\alpha} = B^{ij}$$

Are the RCG and LCG compatibility problems really different?

RCG

$$\frac{\partial y}{\partial x} = F \Leftrightarrow \text{curl } F = 0$$

$$F^T F = \text{specified}(x) \in P_{sym}$$

LCG

$$\frac{\partial y}{\partial x} = F \Leftrightarrow \text{curl } F = 0$$

$$FF^T = B$$

$$\Leftrightarrow F^{-T} F^{-1} = B^{-1}$$

$$\mathfrak{I}(F)^T \mathfrak{I}(F) = \text{specified}(x) \in P_{sym}$$

Using RCG method one has

$$\text{curl } \mathfrak{I}(F(x)) = 0$$

$$\varepsilon_{ijk} \frac{\partial F_{an}}{\partial x^j} F_{nk}^{-1} = 0 \text{ but need } \varepsilon_{ijk} \frac{\partial F_{ak}}{\partial x^j} = 0$$

Motivation

■ Continuum Mechanics

- Interesting geometry question for classical kinematical measure
 - Cauchy stress for Frame-indifferent, isotropic elastic material is a function only of B
- Sharp contrast in uniqueness from the more studied RCG case
- Open in 3-D

■ Mathematics

- Interesting questions involving
- Geometry
- Nonlinear PDE
- Algebra
 - Even in the C^∞ local existence case

Riemannian Geometry in charts

Given two coordinate patches for a Riemannian manifold, with points denoted generically by

$$x \Leftrightarrow x^\alpha \Leftrightarrow (x^1, x^2, x^3) \quad \text{and} \quad y = y(x)$$

$$y \Leftrightarrow y^\alpha \Leftrightarrow (y^1, y^2, y^3) \quad \text{and} \quad \det\left(\frac{\partial y}{\partial x}\right) \neq 0$$

\exists 3×3 sym, + def matrix fields

$${}_{(x)}C, \quad {}_{(y)}C, \quad {}_{(x)}B := {}_{(x)}C^{-1}, \quad {}_{(y)}B := {}_{(y)}C^{-1}$$

satisfying

$${}_{(x)}C_{\alpha\beta} = \frac{\partial y^k}{\partial x^\alpha} \left[{}_{(y)}C \right]_{km} \frac{\partial y^m}{\partial x^\beta}$$

Matrix
inverse

$${}_{(y)}B^{ij} = \frac{\partial y^i}{\partial x^\alpha} \left[{}_{(x)}B \right]^{\alpha\beta} \frac{\partial y^j}{\partial x^\beta}$$

$$\left[\partial y / \partial x \right]_{k\alpha} = \frac{\partial y^k}{\partial x^\alpha}$$

Notational agreement:
Evaluate anything like

$${}_{(x)}(\cdot) \quad \text{at} \quad x$$

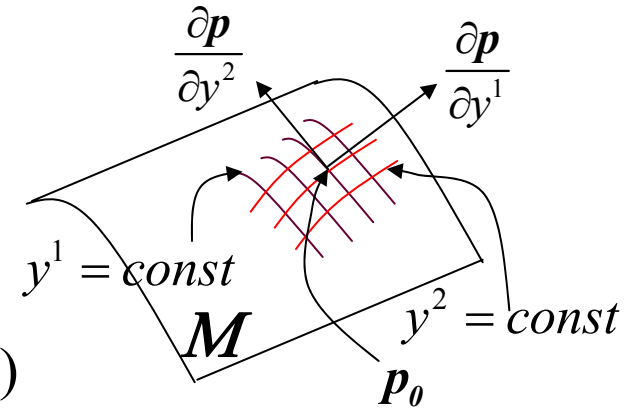
and

$${}_{(y)}(\cdot) \quad \text{at} \quad y(x)$$

The question of equivalence of quadratic forms

Given two local P_{sym} matrix fields on manifold M ,
 $\dim(M) = N$, they are equivalent if one can find
 two local charts in \mathbb{R}^N related 1-1 with

$\det\left(\frac{\partial y}{\partial x}\right) \neq 0$ satisfying transformation rules. (WHY???)



On tangent space T_{p_0} of $p_0 \in M$ spanned by

$\left(\frac{\partial p}{\partial y^1}, \frac{\partial p}{\partial y^2}\right)$ or $\left(\frac{\partial p}{\partial x^1}, \frac{\partial p}{\partial x^2}\right)$ any vector $\mathbf{a} \in T_{p_0}$

$$\mathbf{a} = {}_{(x)}a^\alpha \frac{\partial p}{\partial x^\alpha} = {}_{(x)}a^\alpha \frac{\partial p}{\partial y^i} \frac{\partial y^i}{\partial x^\alpha} = {}_{(y)}a^i \frac{\partial p}{\partial y^i} \Rightarrow {}_{(y)}a^i = {}_{(x)}a^\alpha \frac{\partial y^i}{\partial x^\alpha} \left[\frac{\partial y^i}{\partial x^\alpha} {}_{(y)}C_{ij} \frac{\partial y^j}{\partial x^\beta} - {}_{(x)}C_{\alpha\beta} \right] = 0$$

$\forall {}_{(x)}a^\alpha, {}_{(x)}b^\beta$

Now, let there be a quadratic form for each chart s.t. for \mathbf{a}, \mathbf{b}

${}_{(x)}a^\alpha {}_{(x)}C_{\alpha\beta} {}_{(x)}b^\beta =:$ physical scalar indep. of chart

But, no chart is special; $\therefore {}_{(x)}a^\alpha {}_{(x)}C_{\alpha\beta} {}_{(x)}b^\beta = {}_{(y)}a^i {}_{(y)}C_{ij} {}_{(y)}b^j$

Mapping compatibility question to Riemannian Geometry

$${}_{(x)}C_{\alpha\beta} = \frac{\partial y^k}{\partial x^\alpha} \left[{}_{(y)}C \right]_{km} \frac{\partial y^m}{\partial x^\beta}$$

$${}_{(y)}B^{ij} = \frac{\partial y^i}{\partial x^\alpha} \left[{}_{(x)}B \right]^{\alpha\beta} \frac{\partial y^j}{\partial x^\beta}$$

choose

$${}_{(y)}C \equiv I$$

RCG compatibility

$$\frac{\partial y^k}{\partial x^\alpha} \frac{\partial y^k}{\partial x^\beta} = C_{\alpha\beta}$$

given C, B as fns. of x
find $y(x)$

$${}_{(x)}B \equiv I$$

LCG compatibility

$$\frac{\partial y^i}{\partial x^\alpha} \frac{\partial y^j}{\partial x^\alpha} = B^{ij}$$

Overdetermined Problems

Machinery of Riemannian Geometry thanks to Christoffel

Recall notation, for any patch Z $(z)B = (z)C^{-1}$

$${}_{(z)}\Gamma_{rs}^i := \frac{{}_{(z)}B^{ip}}{2} \left[\frac{\partial {}_{(z)}C_{rp}}{\partial z^s} + \frac{\partial {}_{(z)}C_{sp}}{\partial z^r} - \frac{\partial {}_{(z)}C_{rs}}{\partial z^p} \right]$$

Necessary condition for existence of $y(x)$ satisfying metric transformation rules is

$$\frac{\partial y^i}{\partial x^\alpha \partial x^\beta} = {}_{(x)}\Gamma_{\alpha\beta}^\rho \frac{\partial y^i}{\partial x^\rho} - {}_{(y)}\Gamma_{rs}^i \frac{\partial y^r}{\partial x^\alpha} \frac{\partial y^s}{\partial x^\beta}$$

Roughly:

For RCG ${}_{(y)}C \equiv I \Rightarrow {}_{(y)}\Gamma_{rs}^i \equiv 0$ Linear problem for $\partial y / \partial x$

For LCG ${}_{(x)}B \equiv I \Rightarrow {}_{(x)}\Gamma_{\alpha\beta}^\rho \equiv 0$ Quasilinear problem for $\partial y / \partial x$!!!!

RCG compatibility

- Riemann
- Christoffel
- Brothers Cosserat
-
-
- Shield
- Deturck and Yang
- Interesting associated facts, especially uniqueness question
 - if two deformations have same RCG field, then they differ by rigid deformation
 - Reshetnyak (according to Ball and James)
 - Inadequacy of single-well energy for prediction of microstructure with compatible elastic deformation (Ball & James)
 - Friesecke, Muller, James
 - An invertible tensor field F may have nonvanishing curl even if its RCG field ($F^T F$) is compatible

Complete Integrability of Pfaff PDE (T.Y. Thomas, 1934)

Theorem: Consider PDE

$$\frac{\partial w^i}{\partial x^\alpha}(x) = \psi_\alpha^i(w(x), x) \quad i = 1 \text{ to } R ; \alpha = 1 \text{ to } n$$

$\psi_\alpha^i \in C^1(\Omega)$, Ω open connected subset of $\mathbb{R}^R \times \mathbb{R}^n$

Suppose the integrability condition

$$\frac{\partial \psi_\alpha^i}{\partial w^j} \psi_\beta^j + \frac{\partial \psi_\alpha^i}{\partial x^\beta} = \frac{\partial \psi_\beta^i}{\partial w^j} \psi_\alpha^j + \frac{\partial \psi_\beta^i}{\partial x^\alpha}$$

holds in Ω . (motivated by equality of second partial derivs.)

Then for arbitrary $(w_0, x_0) \in \Omega$, \exists a unique local solution around x_0 satisfying $w(x_0) = w_0$. Therefore, solution allows R arbitrary constants to be specified.

RCG compatibility

Motivated by necessary condition, consider

$$\frac{\partial y^i}{\partial x^\alpha} = u_\alpha^i$$

$$\frac{\partial u_\alpha^i}{\partial x^\beta} = {}_{(x)}\Gamma_{\alpha\beta}^\gamma u_\gamma^i = u_\gamma^i \frac{C^{\gamma\mu}}{2} \left[\frac{\partial C_{\alpha\mu}}{\partial x^\beta} + \frac{\partial C_{\beta\mu}}{\partial x^\alpha} - \frac{\partial C_{\alpha\beta}}{\partial x^\mu} \right]$$

${}_{(x)}B$

Integrability Condition (nice and 'separably' factored in x -dependent terms and u -dependent terms)

$$u_\mu^i \left[\frac{\partial {}_{(x)}\Gamma_{\alpha\beta}^\mu}{\partial x^\rho} - \frac{\partial {}_{(x)}\Gamma_{\alpha\rho}^\mu}{\partial x^\beta} + {}_{(x)}\Gamma_{\gamma\rho}^\mu {}_{(x)}\Gamma_{\alpha\beta}^\gamma - {}_{(x)}\Gamma_{\gamma\beta}^\mu {}_{(x)}\Gamma_{\alpha\rho}^\gamma \right] = 0$$

\therefore require Riemann-Christoffel curvature tensor to vanish. Guarantees existence of u with arbitrarily specifiable value u_0 at one point, and because of symmetry of Γ in lower indices, of y .

Remains to be shown that $u_\alpha^i u_\beta^i = C_{\alpha\beta}$

RCG Compatibility

Assign $u(x_0)$ such that $u^T u(x_0) = C(x_0)$.

Continuity $\Rightarrow u$ invertible locally around x_0 .

Define $v = u^{-1}$; noting $\delta_{ij} = v_i^\alpha C_{\alpha\beta} v_j^\beta(x_0)$

$$\frac{\partial}{\partial x^\mu} (v_i^\alpha C_{\alpha\beta} v_j^\beta) = v_j^\beta v_i^\alpha \left[\frac{\partial C_{\rho\beta}}{\partial x^\mu} - C_{\alpha\beta} \Gamma_{\rho\mu}^\alpha - C_{\rho\alpha} \Gamma_{\beta\mu}^\alpha \right] = 0 \quad \text{!!!!}$$

covariant derivative of covariant metric tensor

Ricci: metric tensors covariantly constant

(merely smoothness, and defn. of Γ !!!)

$$\therefore \delta_{ij} = v_i^\alpha C_{\alpha\beta} v_j^\beta \quad \Rightarrow \quad u^T u = C = \left(\frac{\partial y}{\partial x} \right)^T \frac{\partial y}{\partial x} \quad \text{locally } \square$$

Left Cauchy Green Compatibility

- 2-D
 - **Blume (1989)**
 - Formulation based on Polar Decomposition (for both 2/3-D)
 - Find a rotation tensor field
 - Compatibility condition for 2-d problem is derived
 - ‘Explicit’ characterization of the condition
 - Uniqueness is analyzed
 - **Duda & Martins (1995)**
 - Plane case
 - Polar Decomposition
 - Analysis of possible cases; construction of the rotation field
 - Insightful and detailed analysis of the uniqueness question
 - Demonstration of nonuniqueness through constructive examples systematically using Thomas
- 3-D - open
 - **Acharya (1999)**
 - geometric formulation
 - provides condition for local existence in 3-d
 - Much can be done in ‘explicit’ characterization of the existence condition

LCG Compatibility

Consider necessary condition for existence of y^i such that

$$\frac{\partial y^i}{\partial x^\alpha} \frac{\partial y^j}{\partial x^\alpha} (x) = {}_{(y)}B^{ij} (x) \text{ holds. (and } {}_{(x)}B^{\alpha\beta} (x) = \delta^{\alpha\beta} \text{ constant)}$$

Around arbitrary x_0 , the map y is then locally invertible and so

$$\frac{\partial B_{rp}}{\partial y^s} = \frac{\partial B_{rp}}{\partial x^\alpha} \frac{\partial x^\alpha}{\partial y^s} \quad \text{where } \left[{}_{(y)}B^{-1} \right]_{rp} =: B_{rp} = {}_{(y)}C_{rp}$$

$$\text{Recall } {}_{(z)}B = {}_{(z)}C^{-1} \quad {}_{(z)}\Gamma_{rs}^i := \frac{{}_{(z)}B^{ip}}{2} \left[\frac{\partial {}_{(z)}C_{rp}}{\partial z^s} + \frac{\partial {}_{(z)}C_{sp}}{\partial z^r} - \frac{\partial {}_{(z)}C_{rs}}{\partial z^p} \right]$$

$$\frac{\partial y^i}{\partial x^\alpha \partial x^\beta} = {}_{(x)}\Gamma_{\alpha\beta}^\rho \frac{\partial y^i}{\partial x^\rho} - {}_{(y)}\Gamma_{rs}^i \frac{\partial y^r}{\partial x^\alpha} \frac{\partial y^s}{\partial x^\beta}$$

$$\therefore {}_{(x)}\Gamma_{\alpha\beta}^\rho \equiv 0, \text{ and } \frac{\partial y^i}{\partial x^\alpha \partial x^\beta} = -\frac{B^{im}}{2} \left[\frac{\partial B_{rm}}{\partial x^\beta} \frac{\partial y^r}{\partial x^\alpha} + \frac{\partial B_{sm}}{\partial x^\alpha} \frac{\partial y^s}{\partial x^\beta} - \frac{\partial B_{rs}}{\partial x^\rho} \frac{\partial x^\rho}{\partial y^m} \frac{\partial y^r}{\partial x^\alpha} \frac{\partial y^s}{\partial x^\beta} \right]$$

Left Cauchy Green Compatibility Governing PDE System

Original – nonlinear, first order system

Formulated as – Quasilinear, Pfaffian system

$$u_{\rho}^i u_{\rho}^j = B^{ij}$$

$$\left[{}_{(y)}B \right]^{ij} (x) =: [B]^{ij} (x)$$

$$\left[{}_{(y)}B^{-1} \right]_{rp} =: B_{rp} = {}_{(y)}C_{rp}$$

$$\frac{\partial y^i}{\partial x^{\alpha}} = u_{\alpha}^i$$

Matrix inverse function

$$\frac{\partial u_{\alpha}^i}{\partial x^{\beta}} = -\frac{B^{im}}{2} \left[\frac{\partial B_{rm}}{\partial x^{\beta}} u_{\alpha}^r + \frac{\partial B_{sm}}{\partial x^{\alpha}} u_{\beta}^s - \frac{\partial B_{rs}}{\partial x^{\rho}} \mathfrak{I}_m^{\rho}(u) u_{\alpha}^r u_{\beta}^s \right]$$

Setup

Seek functions

$$w^i(x), \quad i = 1, 2, \dots, R$$

that satisfy

$$\frac{\partial w^i}{\partial x^\alpha}(x) = \psi_\alpha^i(w(x), x) \quad \alpha = 1, 2, \dots, n$$

(for definiteness think of domain of ψ_α^i to be open connected set of $\mathbb{R}^R \times \mathbb{R}^n$)

Refer to domain of ψ_α^i as (z, x)

Associate $w^i \rightarrow u_\alpha^i$; so $R = 3 \times 3 = 9$; $n = 3$.

$$\frac{\partial u_\alpha^i}{\partial x^\beta} = A_{\alpha\beta}^i \left(u, \quad {}_{(y)}B(x), \quad \frac{\partial {}_{(y)}C}{\partial x}(x) \right) ; \quad \frac{\partial y^i}{\partial x^\alpha} = u_\alpha^i$$

Sufficient condition for local existence: the completely-integrable situation

Hypothesis: Suppose

$$F^{(1)}(u, x) := \left(\frac{\partial A_{\alpha\beta}^i}{\partial u_{\mu}^k} A_{\mu\gamma}^k + \frac{\partial A_{\alpha\beta}^i}{\partial x^{\gamma}} - \frac{\partial A_{\alpha\gamma}^i}{\partial u_{\mu}^k} A_{\mu\beta}^k - \frac{\partial A_{\alpha\gamma}^i}{\partial x^{\mu}} \right) (u, x) \equiv 0 \text{ locally in } (u, x)$$

(seek symmetry in β, γ for each $i, \alpha \Rightarrow 27$ nonlinear algebraic equations)

If so, Thomas guarantees solution to u and therefore y with arbitrary data at one x_0 .

So - specify conditions on B field for when identity can be satisfied:

Unlike RCG case, $F^{(1)}$ is

● cumbersome (downright scary!) ● nonlinear in u ● does not readily

SEPARABLY factorize into at least 1 solely x – dependent term

- need separability (seems to me) for identity with control only on x – dependent terms

- need (algebraic-geometric?) theorem for when this can happen

given field ${}_{(y)}B(x)$ and algebraic structure of array A .

Sufficient condition for local existence: the completely-integrable situation

Initial data to match $uu^T(x_0) = {}_{(y)}B(x_0) \in P_{ym}$ can be constructed. Then local diffeomorphism y around x_0 satisfying $\partial y / \partial x = u$ exists (and u is invertible)

Now define ${}_{(y)}\Gamma_{rs}^i(y') := \left(\frac{{}_{(y)}B^{ip} \circ x}{2} \left[\frac{\partial {}_{(y)}C_{rp} \circ x}{\partial y^s} + \frac{\partial {}_{(y)}C_{sp} \circ x}{\partial y^r} - \frac{\partial {}_{(y)}C_{rs} \circ x}{\partial y^p} \right] \right) (y')$

(Notation: $y^{-1} =: x$) Then,

$$\frac{\partial u_\alpha^i}{\partial x^\beta}(x(y')) = -{}_{(y)}\Gamma_{rs}^i(y') [u_\alpha^r u_\beta^s](x(y'))$$

Define $v = u^{-1}$, consider $\frac{\partial}{\partial y^m} \left[(B^{ij} v_i^\alpha v_j^\beta) \circ x \right]$

$$\frac{\partial}{\partial y^m} [B^{ij} v_i^\alpha v_j^\beta] = \left[\frac{\partial {}_{(y)}B^{ij}}{\partial y^m} + {}_{(y)}B^{kj} {}_{(y)}\Gamma_{km}^i + {}_{(y)}B^{ik} {}_{(y)}\Gamma_{km}^j \right] v_i^\alpha v_j^\beta = 0$$

Ricci: covariant deriv. of
contravariant Metric tensor = 0

Since u invertible,

$$\frac{\partial [\]}{\partial x} u^{-1} = 0 \Rightarrow \frac{\partial y}{\partial x} \left(\frac{\partial y}{\partial x} \right)^T = B$$

□

Sufficient condition for NOT completely integrable case

Let $F^{(1)}(z, x) \neq 0$ identically. ($F^{(1)}$ defines complete integrability condition)

Define $F_\alpha^{(j+1)}(z, x) := \left(\sum_{i=1}^R \frac{\partial F^{(j)}}{\partial z^i} \psi_\alpha^i + \frac{\partial F^{(j)}}{\partial x^\alpha} \right)(z, x)$

and consider two integers N, R with

$$1 \leq N \leq R \quad ; \quad 1 \leq M \leq R.$$

Assume that

● there exist M equations in the sets $F^{(1)} = 0$ through $F^{(N)} = 0$ denoted by $\tilde{G}_\lambda = 0$, $\lambda = 1$ to M ,

and M of the variables z^i (from the list $z^i, i = 1$ to R) denoted by

\bar{z}^i , $i = 1$ to M identified through a known one-to-one map $\kappa : \{1, 2, \dots, M\} \rightarrow \{1, 2, \dots, R\}$

by $\bar{z}^i := z^{\kappa(i)}$

which satisfy $\det \left[\frac{\partial \tilde{G}_\lambda}{\partial \bar{z}^i}(z, x) \right] \neq 0$ locally in (z, x) space.

Not Completely Integrable Case, contd.

Denote remaining $R - M =: P$ variables z as $\hat{z}^i, i = 1$ to P , defined by

$$\mu: \{1, 2, \dots, P\} \rightarrow \{1, 2, \dots, R\} \quad \hat{z}^i := z^{\mu(i)}.$$

● Assume that around a point $(\hat{z}_0, x_0) := (\hat{z}_0^1, \dots, \hat{z}_0^P, x_0^1, \dots, x_0^n)$, the solution

$$\bar{z}^i = \varphi^i(\hat{z}_0^1, \dots, \hat{z}_0^P, x_0^1, \dots, x_0^n), \quad i = 1, 2, \dots, M, \text{ of } \tilde{G}_\lambda = 0, \quad \lambda = 1, 2, \dots, M$$

satisfies all the equations of the sets

$$F^{(1)} = 0 \quad \text{through} \quad F^{(N+1)} = 0$$

identically in a neighborhood of (\hat{z}_0, x_0) in (\hat{z}, x) space.

Then \exists a function $w(x)$ satisfying

$$\frac{\partial w^i}{\partial x^\alpha}(x) = \psi_\alpha^i(w(x), x) \quad \alpha = 1, 2, \dots, n, \quad i = 1, 2, \dots, R$$

determined by P constants.

Proof: Not completely Integrable case

Let $\Pi(\bar{z}, \hat{z}, x) \mapsto (z, x)$. Define,

$$G_\lambda(\bar{z}, \hat{z}, x) = \tilde{G}_\lambda(\Pi(\bar{z}, \hat{z}, x)), \quad \lambda = 1, 2, \dots, M$$

$$\bar{\psi}_\alpha^i(\bar{z}, \hat{z}, x) = \psi_\alpha^{\kappa(i)}(\Pi(\bar{z}, \hat{z}, x)), \quad i = 1, 2, \dots, M, \quad \alpha = 1, 2, \dots, n$$

$$\hat{\psi}_\alpha^j(\bar{z}, \hat{z}, x) = \psi_\alpha^{\mu(j)}(\Pi(\bar{z}, \hat{z}, x)), \quad j = 1, 2, \dots, P, \quad \alpha = 1, 2, \dots, n$$

Since $F^{(N+1)}(\varphi(\hat{z}, x), \hat{z}, x) \equiv 0$ in (\hat{z}, x)

$$\left[\frac{\partial G_\lambda}{\partial \bar{z}^i} \bar{\psi}_\alpha^i + \frac{\partial G_\lambda}{\partial \hat{z}^j} \hat{\psi}_\alpha^j + \frac{\partial G_\lambda}{\partial x^\alpha} \right] (\varphi(\hat{z}, x), \hat{z}, x) = 0 \quad (I)$$

For arbitrarily fixed path $f(t)$ with $f(0) = x^* \exists$ solution to ODE

$$\frac{dg^j}{dt}(t) = \left[\hat{\psi}_\alpha^j \circ (\varphi \circ (g, f), g, f) \frac{df^\alpha}{dt} \right] (t) \quad ; \quad g^j(0) = \hat{z}^* \text{ (arbitrary),}$$

and g satisfies $G_\lambda \circ (\varphi \circ (g, f), g, f)(t) = 0$ for t in some local interval around 0.

Proof: Not Completely Integrable Case

Differentiating w.r.t. t , combining with (I), and using arbitrariness of path f , \hat{z}^* , and x^* yields,

$$\bar{\psi}_\alpha^i(\varphi(\hat{z}, x), \hat{z}, x) - \frac{\partial \varphi^i}{\partial \hat{z}^j}(\hat{z}, x) \hat{\psi}_\alpha^j(\varphi(\hat{z}, x), \hat{z}, x) - \frac{\partial \varphi^i}{\partial x^\alpha}(\hat{z}, x) \equiv 0 \quad (II)$$

locally in (\hat{z}, x) space.

$F^{(1)}(\varphi(\hat{z}, x), \hat{z}, x) \equiv 0$ combined with (II) implies that the conditions of complete integrability for

$$\frac{\partial g^j}{\partial x^\alpha}(x) = \hat{\psi}_\alpha^j(\varphi(g(x), x), g(x), x) \quad (III)$$

are satisfied. Thus \exists solution to (III) with P arbitrary constants.

Then it can be shown that

$$h^i(x) := \varphi^i(g(x), x) \quad i = 1, 2, \dots, M$$

$$g^j(x) \quad j = 1, 2, \dots, P$$

satisfy $\frac{\partial h^i}{\partial x^\alpha}(x) = \psi^{\kappa(i)} \circ (h, g, Id)(x) \quad \frac{\partial g^j}{\partial x^\alpha}(x) = \psi^{\mu(j)} \circ (h, g, Id)(x). \quad \square$

If $u_0 u_0^T = B(x_0)$ cannot be accommodated by P constants, stick in $uu^T(x) = B(x)$ in $F^{(1)}$. Then use standard idea with Ricci to show B compatibility.

Comments

- Proof adapted from Eisenhart, 1927, Veblen and Thomas, 1926
 - They claim necessity as well; I couldn't
- Guess:
 - Related to Cartan's Method of Equivalence
 - Modern treatment – R B Gardner CBMS-NSF-SIAM