

**In celebration of *Fong Shih*  
(American Academy of Arts and Sciences,  
Cambridge, MA, 12 May 2016)**

***Mechanics on our Planet:  
Great Ice Sheets, Rapid Earthquakes***

***James R. Rice (Harvard) and Ares J. Rosakis (Caltech)***

***April 2016***

# What's a glacier?

- “A river of ice” that flows under its own weight
- Different types:
  - Mountain (valley), Tidewater, Ice sheet

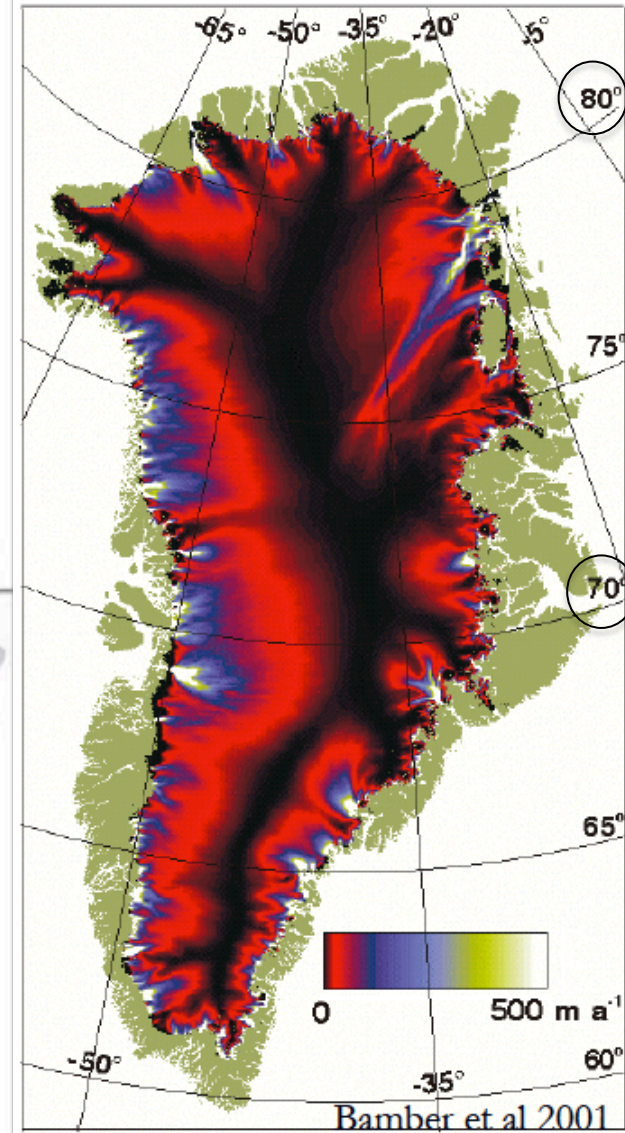
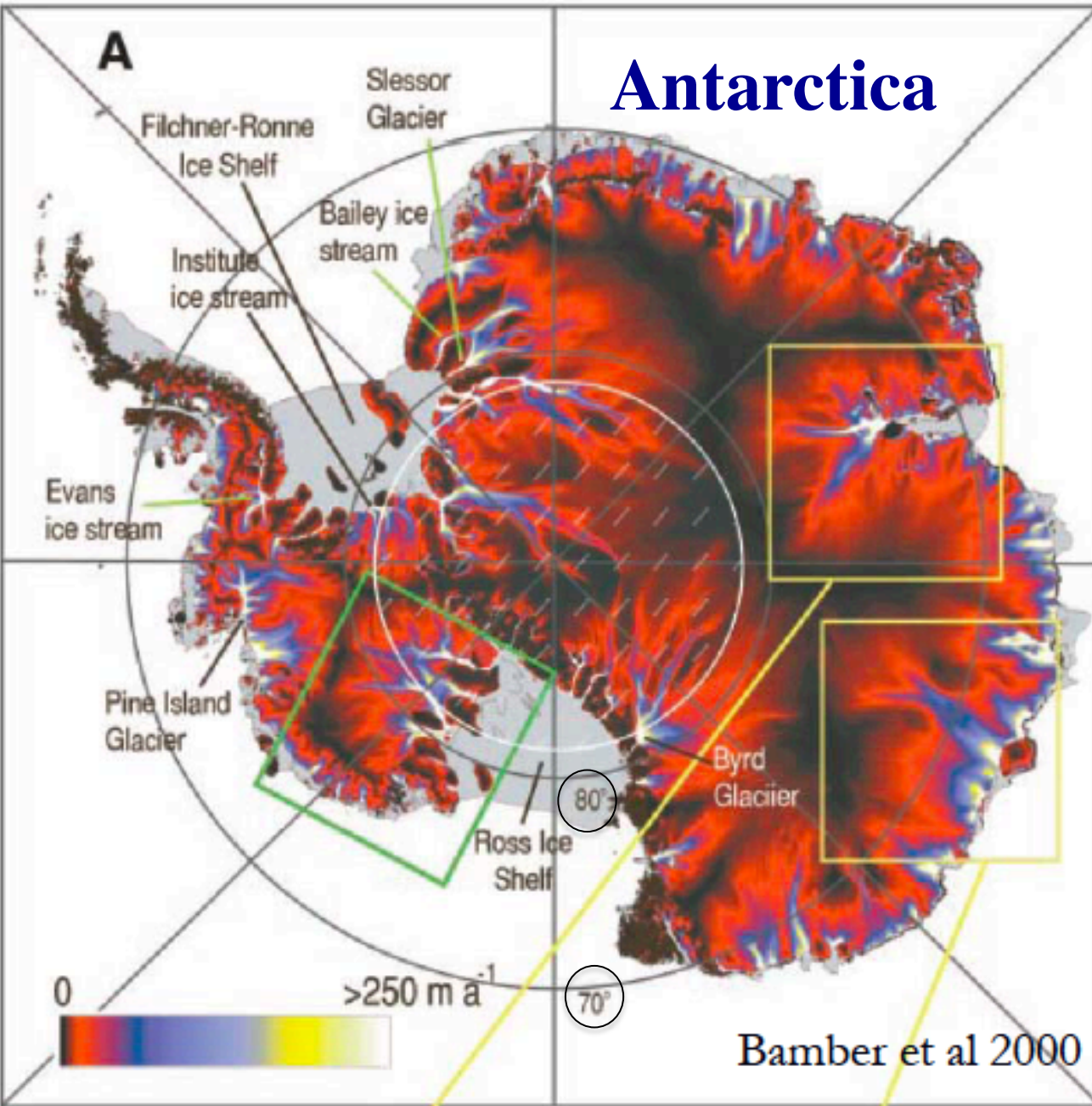


Described in 11<sup>th</sup> century

Motion not realized until 16<sup>th</sup> century!!

# The major ice sheets - *not to scale (reduce Greenland ~50%)*

## Greenland



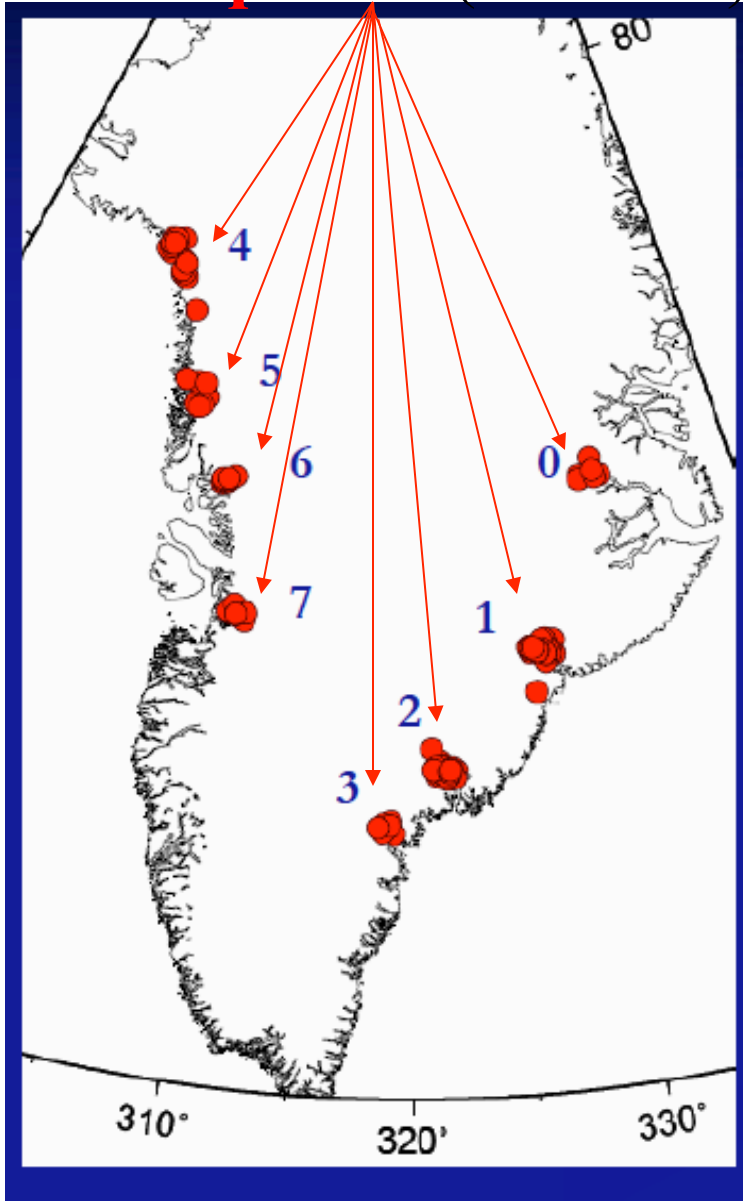
*On the Greenland Ice Sheet:*

*Glacial Earthquakes (discovered and  
located by Göran Ekström), and  
their unexpected mechanism*

*with*

*Victor C. Tsai (Caltech) and, later,  
Mark Fahnstock (Univ. New Hampshire)*

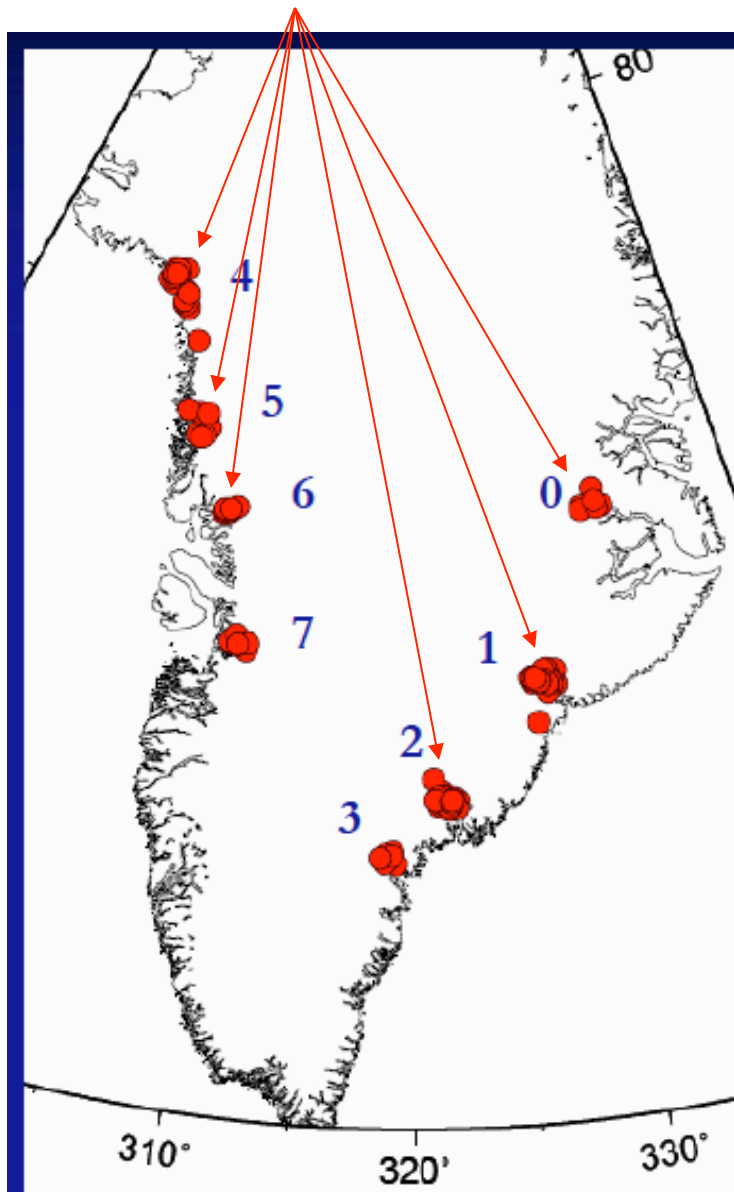
## Source locations of **glacial earthquakes** (Ekström)



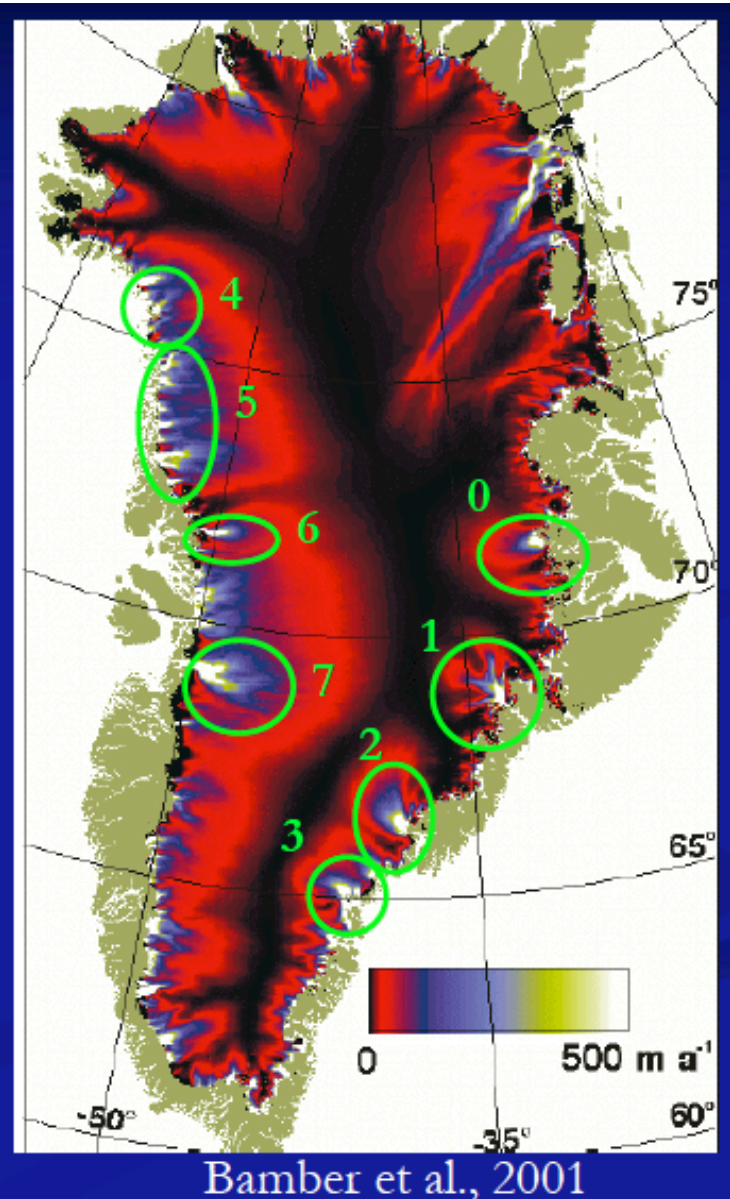
### *Unusual earthquakes:*

- Magnitude  $M_{\text{sw}} \sim 4.6$  to 5.1, measured from signals filtered to 35-150 sec periods.
- Significant energy in periods between 20 and 100 sec (**much longer period than for standard earthquakes of similar  $M_s$**  (e.g., typical source duration  $\sim 2$  sec for  $M_s \sim 5$ )).

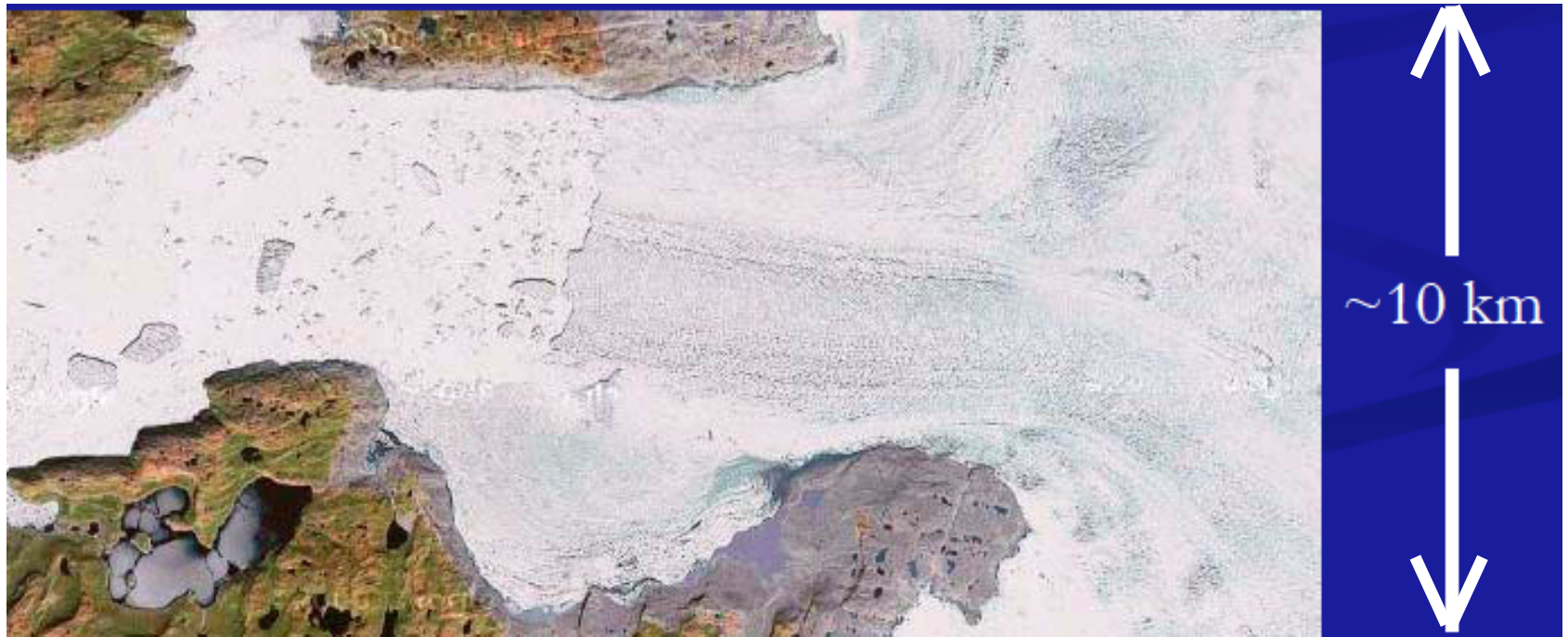
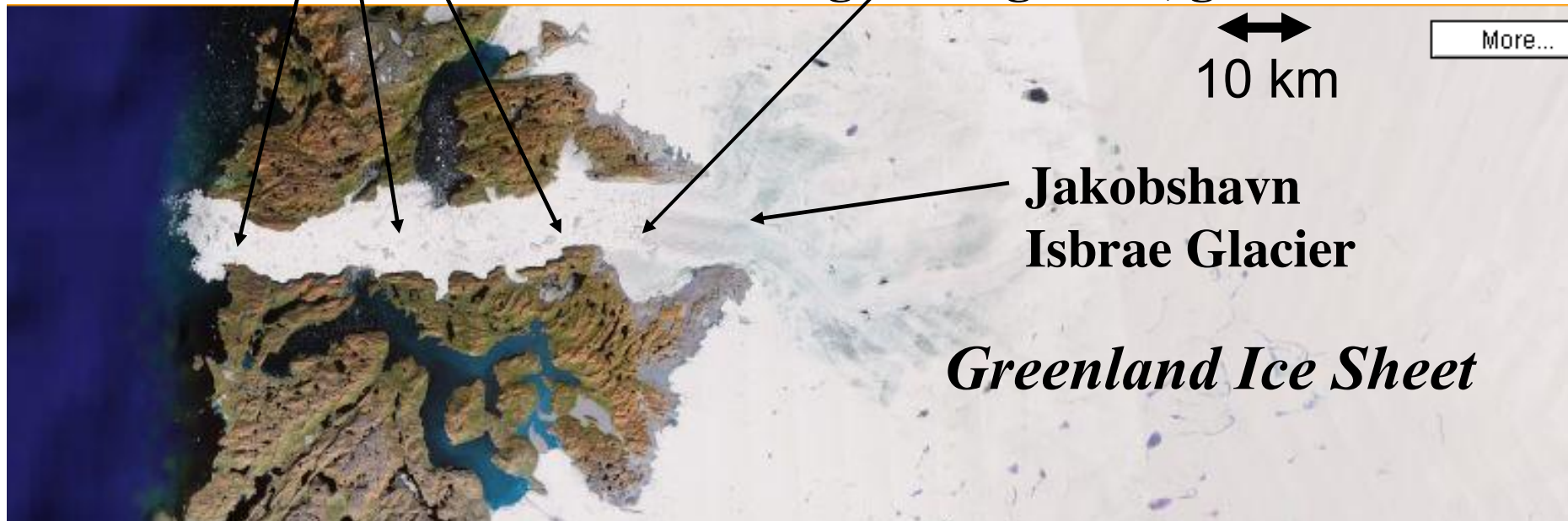
# Source locations of **glacial earthquakes** (G. Ekström)



Correlation with areas of high ice flow rates -- at major fjords

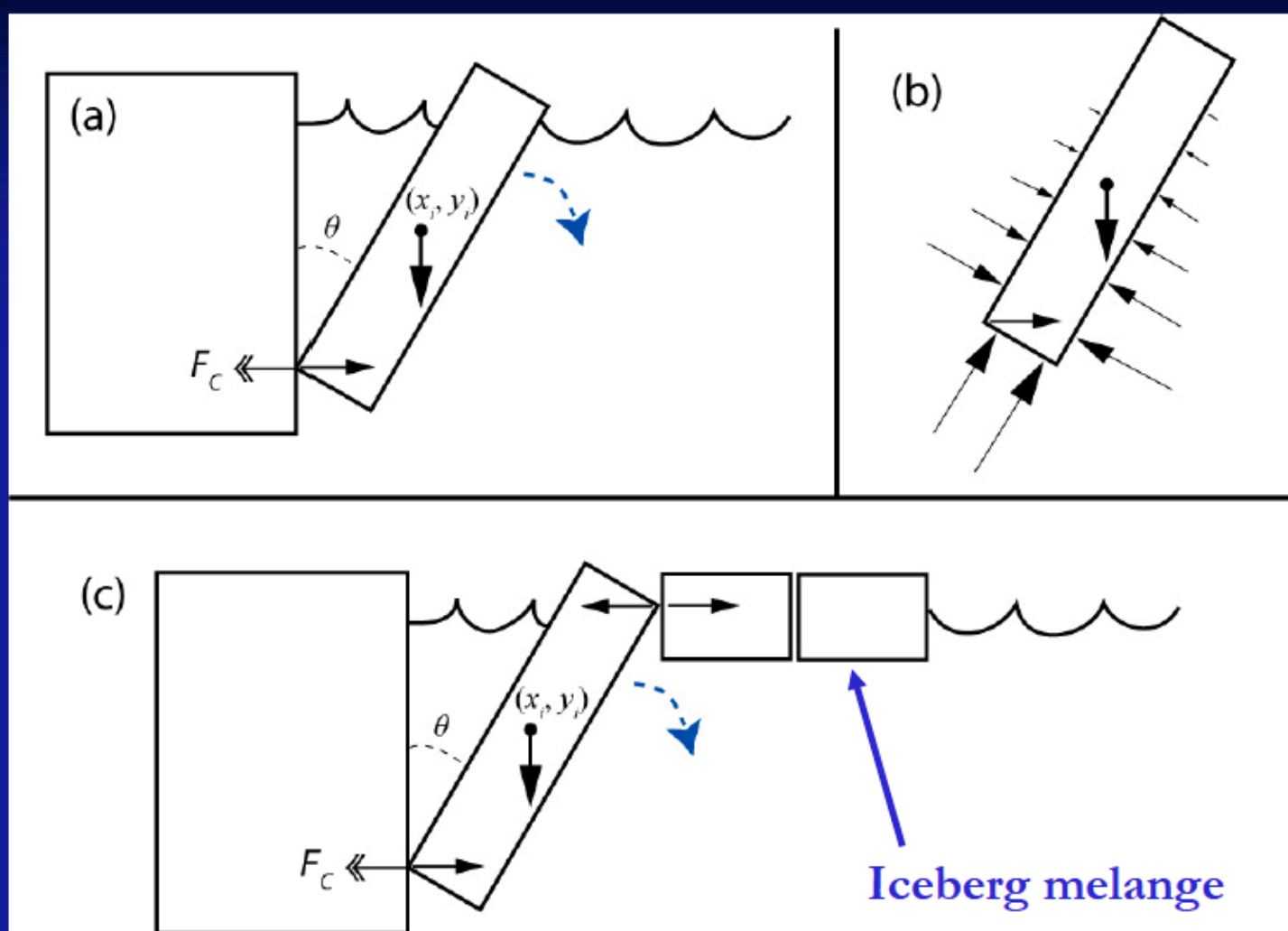


# Melange of calved icebergs      Iceberg calving front, glacier terminus



## What causes glacial EQs?

- *Fast sliding along boundary of ice sheet, or along internal fault, analogous to normal EQs? Nettles et. al. GPS [AGU,'07]: No!*
- *Simple iceberg calving models work best! -- timescale.*



<http://onlinelibrary.wiley.com/doi/10.1029/2008GL035281/abstract>

Video by Jason Amundson (Univ. Alaska, Fairbanks) is sped-up; large overturns take time of order 1-2 minute. See Amundson et al., GRL 2008)

*Greenland Ice Sheet:*

*Rapidly draining of surficial lakes  
and natural hydraulic fractures*

Studies with

*Victor C. Tsai* (Caltech)

*John D. Platt* (Carnegie Inst. of Sci.)

*Matheus C. Fernandes* (Harvard)

**Melange of calved icebergs**

**Iceberg calving front, at glacier terminus**

**Jakobshavn Isbrae Glacier**

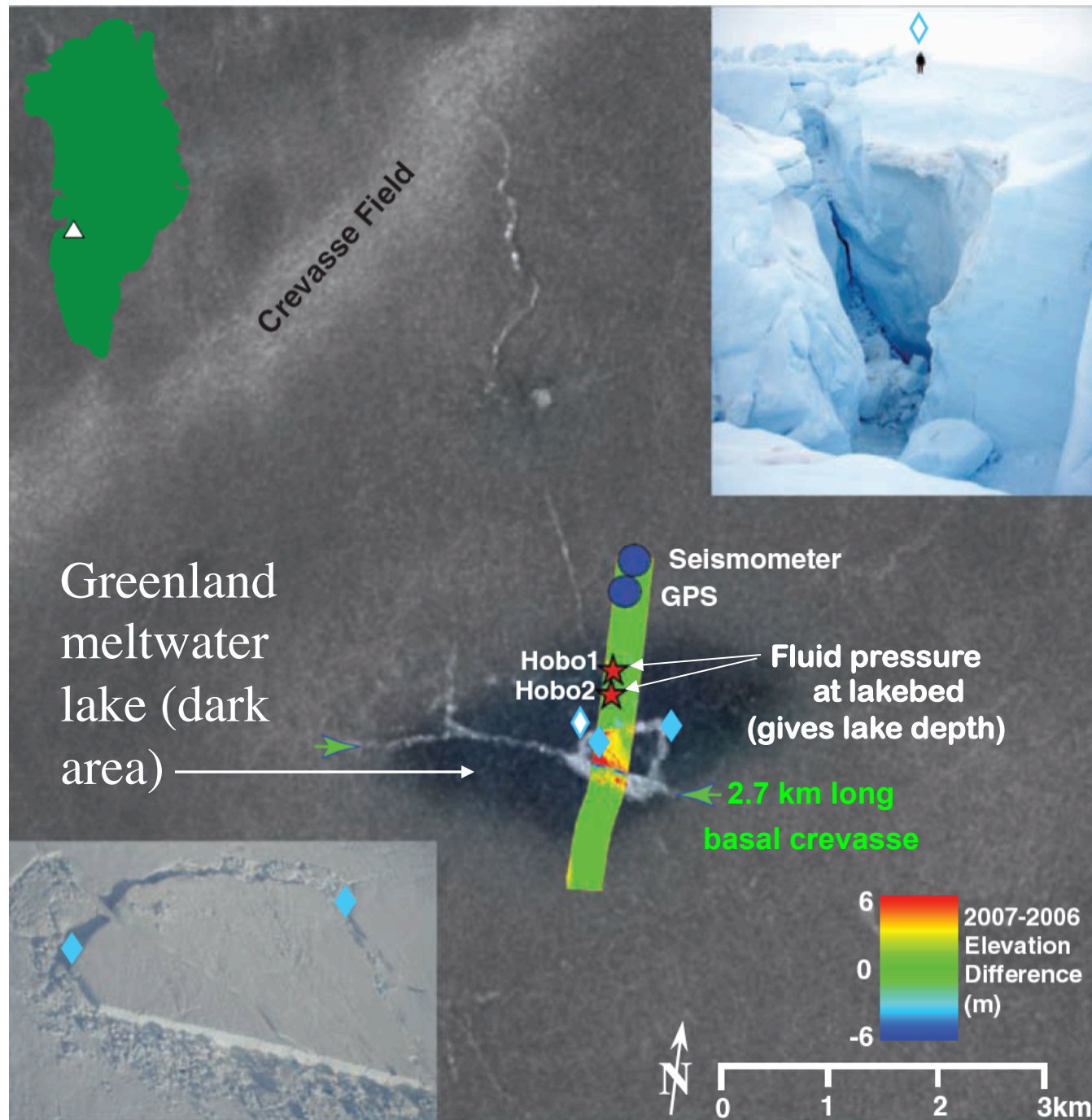
***Greenland Ice Sheet***

**10 km**

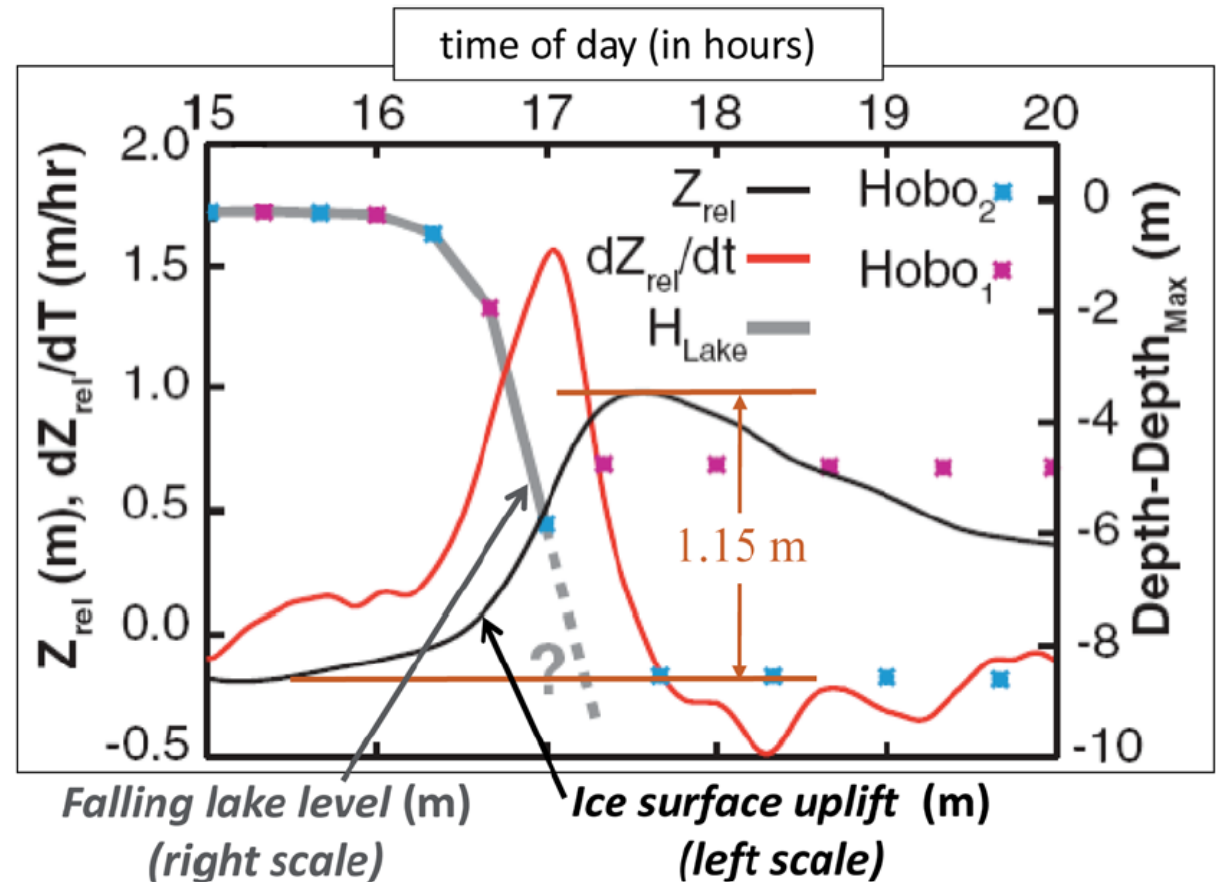
Study motivated by the paper *Fracture Propagation to the Base of the Greenland Ice Sheet During Supraglacial Lake Drainage*, by Das, Joughin, Behn, Howat, King, Lizarralde & Bhatia, *Science*, May 2008.

(Das et al.,  
*Sci.*, 2008)

Early October 2006  
SAR image (gray-  
scale background)  
overlaid with a  
semi-transparent  
image recorded by  
NASA's Moderate  
Resolution Imaging  
Spectroradiometer  
(MODIS) showing  
the lake extent  
(blue) on 29 July  
2006.



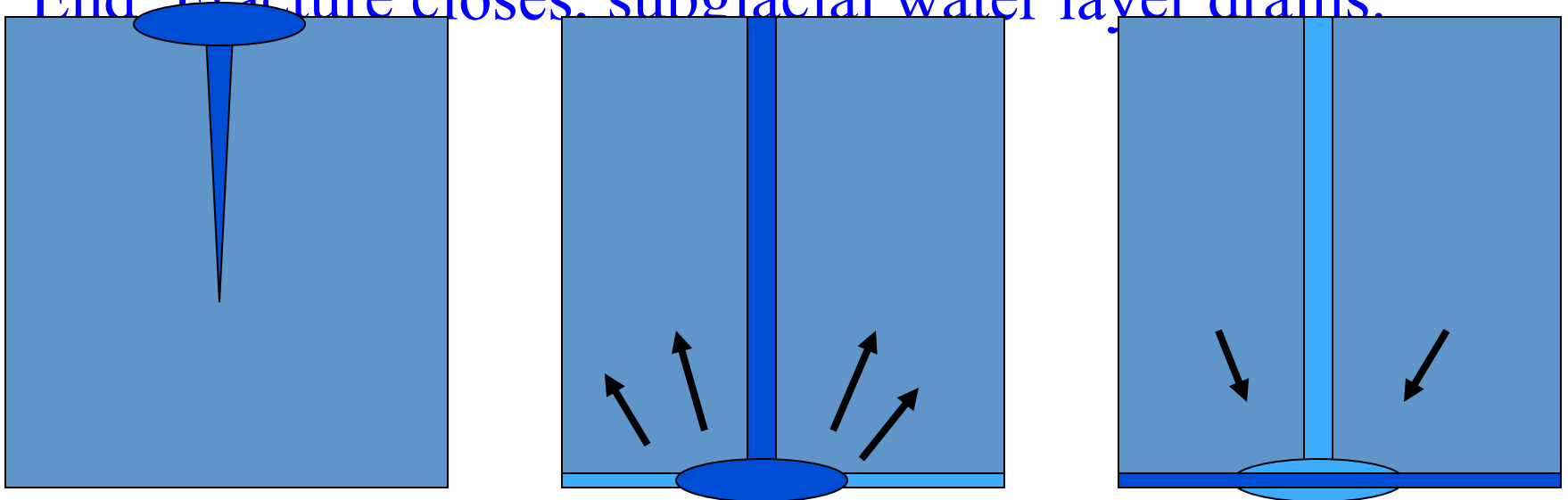
- *Supraglacial meltwater lake* began filling July 2006
- Maximum fill at ~0:00 29 July, Vol. =  $44 \times 10^6 \text{ m}^3$ , Surf. =  $5.6 \text{ km}^2$
- Level slowly/steadily falls,  $\sim 15 \text{ mm/hr}$ , for next  $\sim 16$  hours
- Rapid from 16:00-17:30, max  $12 \text{ m/hr}$  ( $Q_{\text{max}} > 10,000 \text{ m}^3/\text{s}$ ),  
avg  $Q_{\text{avg}} \sim 8,700 \text{ m}^3/\text{s}$  [Compare, Niagra Falls  $Q_{\text{Niag}} \sim 6,000 \text{ m}^3/\text{s}$ ]



[1] Das, Joughin, Behn, Howat, King, Lizarralde & Bhatia, *“Fracture propagation to the base of the Greenland Ice Sheet during supraglacial lake drainage”*, Science, v. 320, pp. 778–781, 2008.

## *Interpretation*

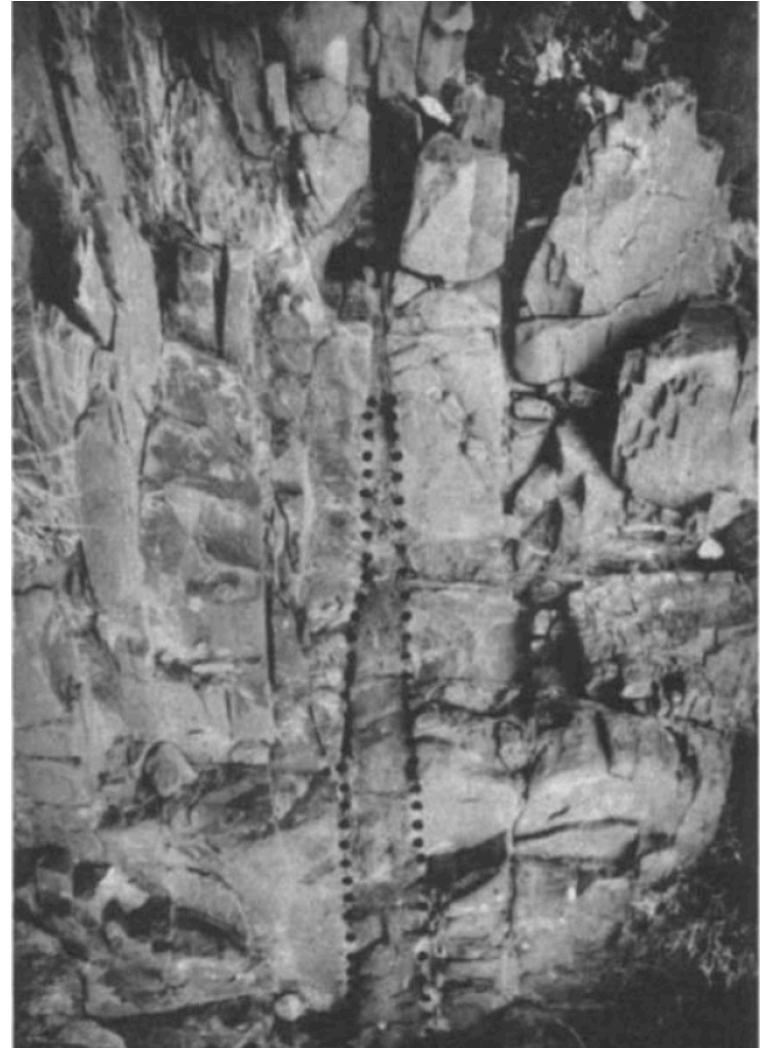
- Initially: Crevasse/moulin system gradually propagates to bed by *Weertman* gravitational instability,  $\rho_{water} > \rho_{ice}$ .
- Middle Stage: Hydraulic cracking and flooding along bed by over-pressure,  $p > \bar{p}_o$  ( $\bar{p}_o$  = ice overburden pressure).
- End: Fracture closes, subglacial water layer drains.



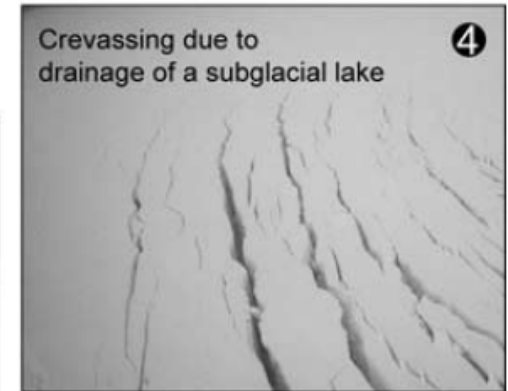
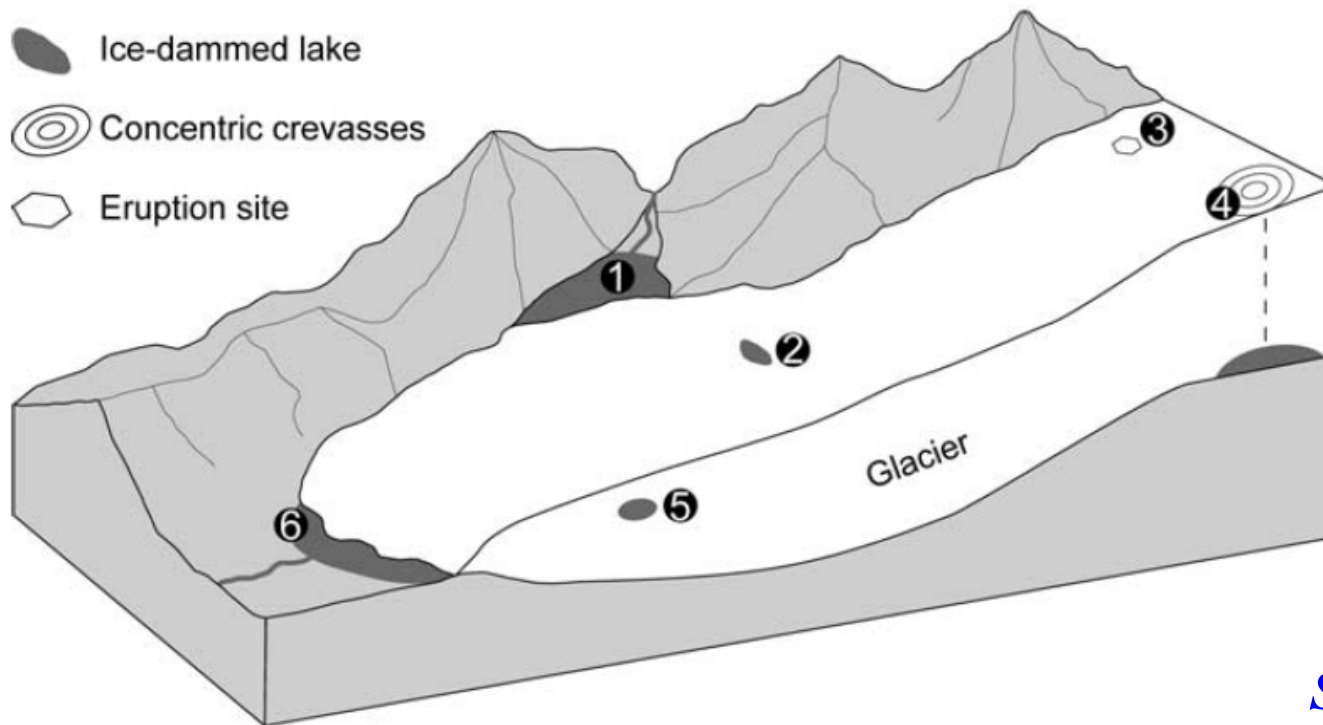
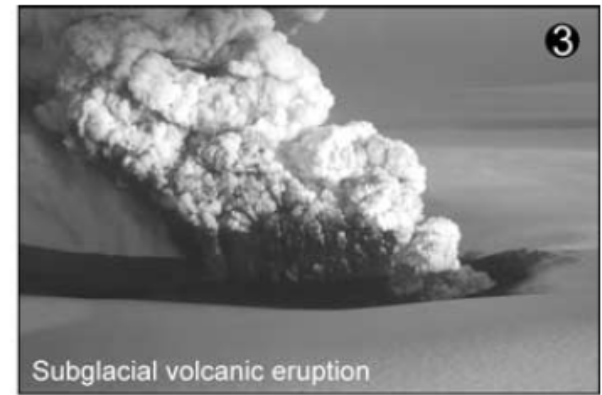
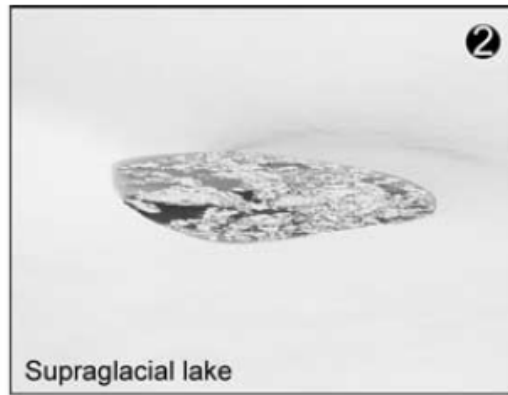
$$\rho_{\text{magma}} < \rho_{\text{rock}}$$



Basaltic dike at tip of Reykjanes Peninsula, southwest Iceland, exposed by glacial erosion (did not make it to surface). Thickness = 40 cm.



Dike (boundaries dotted) terminating in shear zone on Colorado Plateau.



## Other scenarios: *Sub-Glacial Flooding (Jökulhlaup)*

**Figure 1.** Reservoir sites and meltwater sources for jökulhlaups.

from: Roberts, M. J. (2005), Jökulhlaups: A reassessment of floodwater flow through glaciers, *Rev. Geophys.*, 43, RG1002.

# Water drainage along glacier beds often creates channels

Natural example, shown in photo by Laura Kehrl: Rothlisberger Channel  
(but in a mountain glacier, with channel probably created by outburst underflooding):

Kennicott Glacier, Alaska



# Channelized drainage at glacier's bed (Alberta, Canada)



# Alaska Looks for Answers in Glacier's Summer Flood Surges

**New York Times**  
**July 22, 2013**



**First observed  
in July 2011,  
also in July  
2012 & 2013.**

**“... in July  
2011, ... an  
estimated ten  
billion gallons  
gushed out in  
three days ...  
two smaller  
bursts this  
year ...”**

Mathew Ryan Williams for The New York Times

**Visitors leaving caves under the Mendenhall Glacier, near Juneau,  
Alaska. Unpredictable flood surges have elevated concerns.**

By [\*\*KIRK JOHNSON\*\*](#)

**JUNEAU, Alaska -- .... unpredictable flood surges at the [\*\*Mendenhall Glacier\*\*](#),  
about 14 miles from downtown Juneau, Alaska's capital ....**

$$p_{hydrostat} \equiv \rho_{water} g H \geq p_{inlet} \geq \rho_{ice} g H \equiv \sigma_o$$

Rice, Tsai, Fernandes & Platt, *J.Appl.Mech.*, 2015

Average vertical conduit opening:

$$\Delta \bar{u} = \Delta \bar{u}^{el} + \Delta \bar{u}^{cr} \text{ (elastic + prior creep)}$$

$$\Delta \bar{u}^{el} \propto p_{inlet} - \sigma_o$$

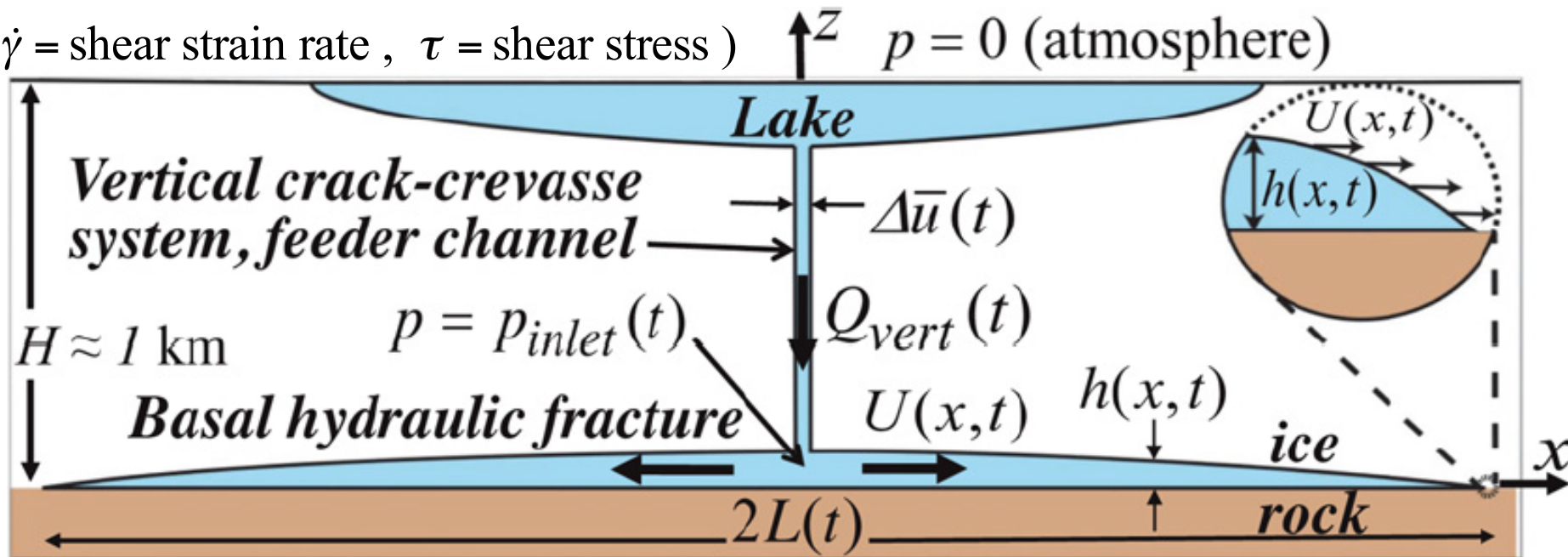
**Important**

For a given conduit opening  $\Delta \bar{u}$ ,  
vertical flow rate

$$Q_{vert} \propto (p_{hydrostat} - p_{inlet})^{1/2}$$

Glen's law, ice deformation:  $\dot{\gamma} = A(T) \tau^3$

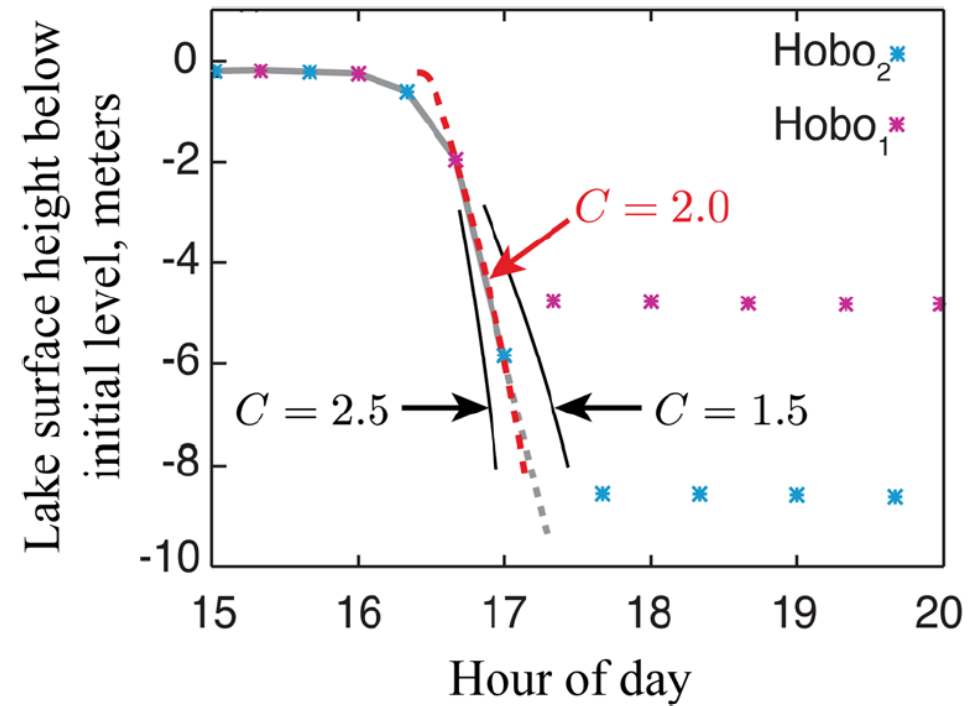
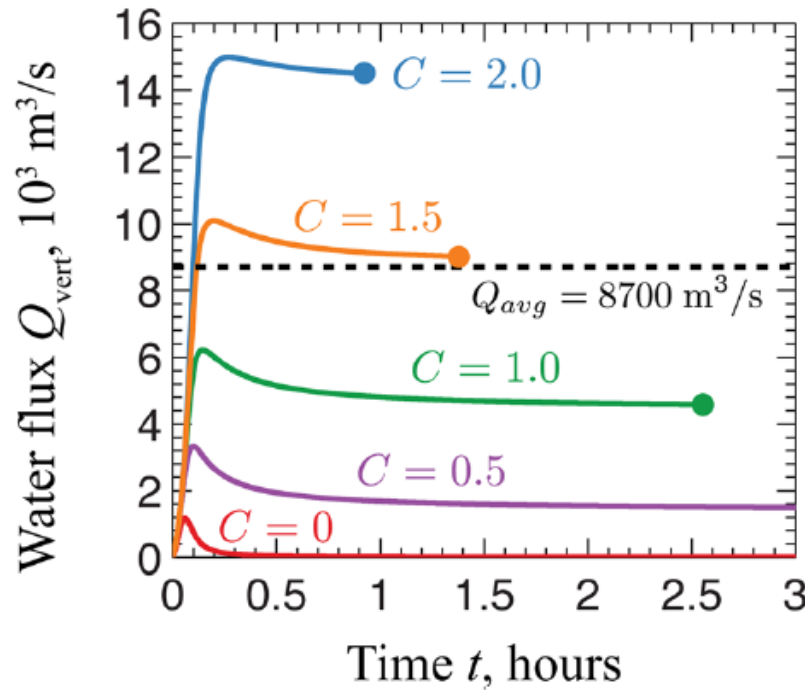
(  $\dot{\gamma}$  = shear strain rate ,  $\tau$  = shear stress )



$p_{inlet} - \sigma_o \sim$  controls flow rate  
 $Q_{basal}$  into basal fracture

$p_{inlet}$  ultimately  
determined by setting  
 $Q_{basal} = Q_{vert}$

$$C = \frac{\text{creep opening over 16 hrs of hydrostatic } p}{\text{elastic opening under hydrostatic } p}$$



At  $T_{\text{ice}} \approx -7^\circ\text{C}$  to  $-5^\circ\text{C}$ , hydrostatic  $p$  over 16 hr of slow leakage  $\Rightarrow C \approx 1.42$  to  $2.11$ , in reasonable agreement with the values of  $C \approx 1.5$  to  $2.0$  which would plausibly fit the the model to the observations. (But, at present, we have no temperature measurements below a persistent supraglacial lake to know if  $-7^\circ\text{C}$  to  $-5^\circ\text{C}$  is sensible.)

Possibility : High stress levels in ice, due to pressurization of the vertical fissure, cause more rapid ice deformation, at any given  $T_{\text{ice}}$ , than predicted by Glenn's law, so that a lower  $T_{\text{ice}}$  would be implied to produce the same drainage time scale.

*West Antarctic Ice Sheet:*

*Rapidly flowing ice streams: What  
processes control their width?*

with

*Thibaut Perol* (Harvard),

*John D. Platt* (Carnegie Inst. Of Sci., DC)

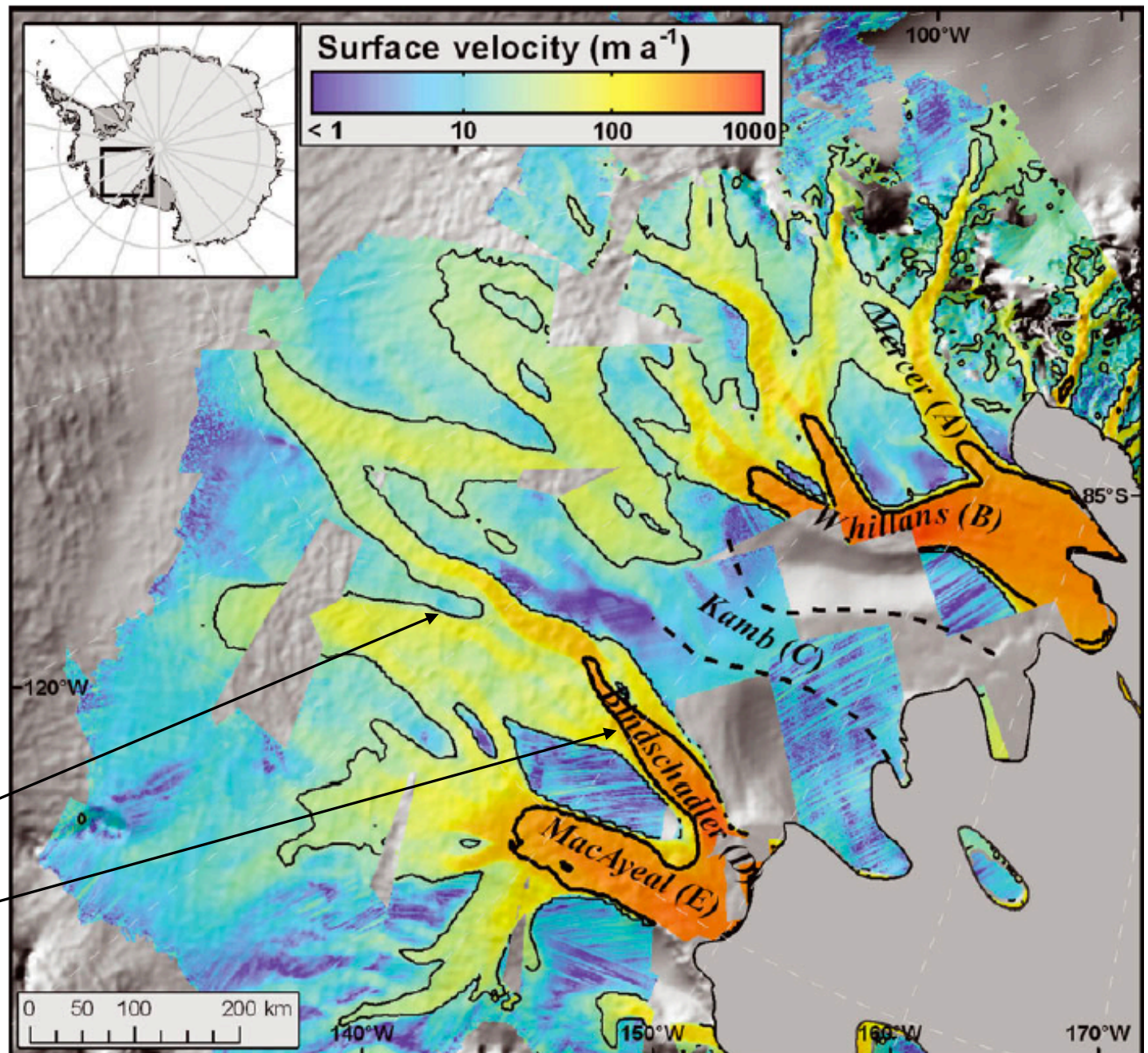
and

*Jenny Suckale* (Stanford)

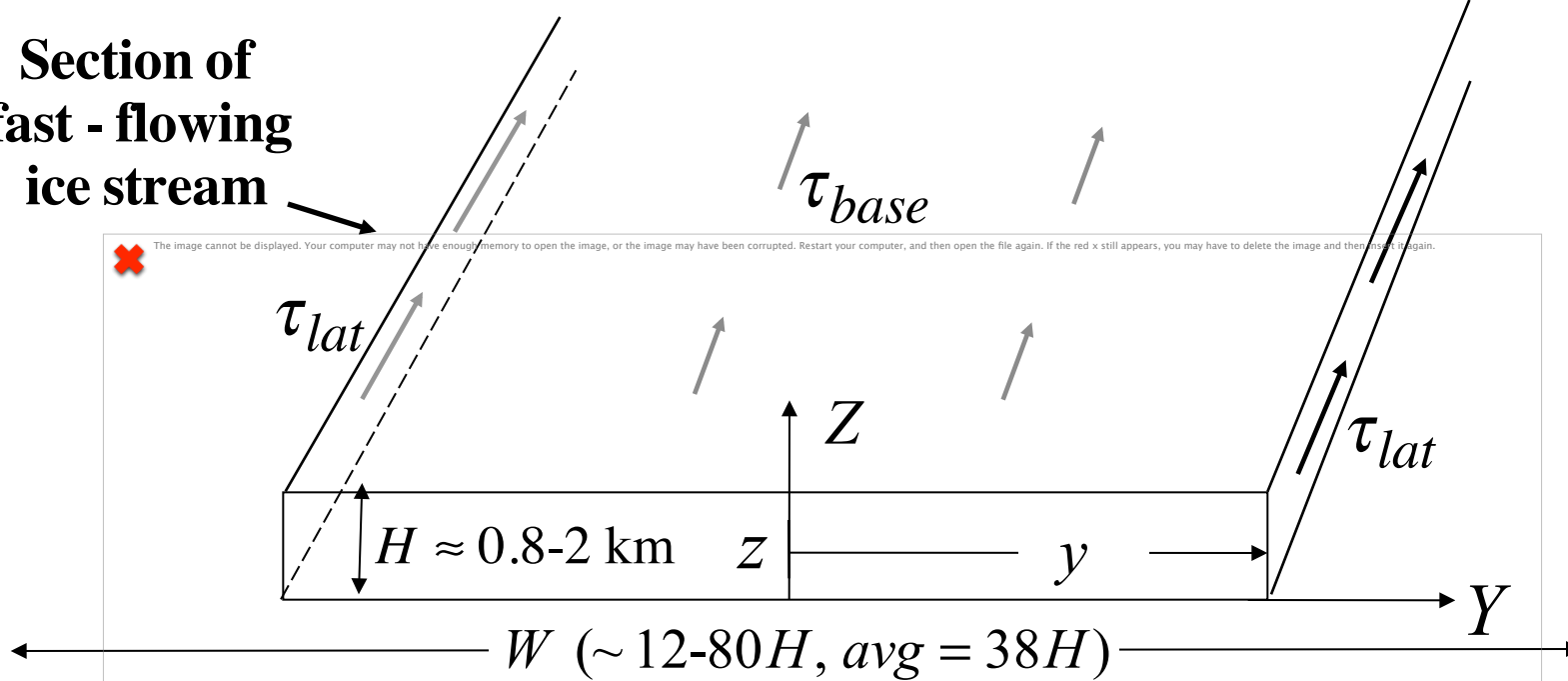
- **Western Antarctica, Siple Coast, Ice Streams,** flowing to the Ross Sea ice shelf.

- InSAR velocity (from Joughin et al., *J. Geoph. Res.*, 2002) overlaid on a digital elevation model (Bamber et al., 2009).

- Velocity contours shown are 25 m/yr (thin line) and 250 m/yr (thick line).



# Section of fast-flowing ice stream



$$\tau_{grav} = \rho g H S \quad (S = \text{slope})$$

= downslope gravity force per unit base area

Equilibrium\*:  $\bar{\tau}_{lat} H = (\tau_{grav} - (\tau_{base})_{avg}) y$

Increases with stream width  $W / H$  ( $\sim 6-40$  at current margins)

$$\bar{\tau}_{lat} \rightarrow (\tau_{grav} - (\tau_{base})_{avg}) \quad \text{low} \quad W / 2H \text{ at margins}$$

⇒ increasing strain rate,

$$\dot{\gamma}_{lat} \sim \bar{\tau}_{lat}^3 \sim (W / H)^3,$$

⇒ increasing shear

$$\text{heating, } \sim (W / H)^4,$$

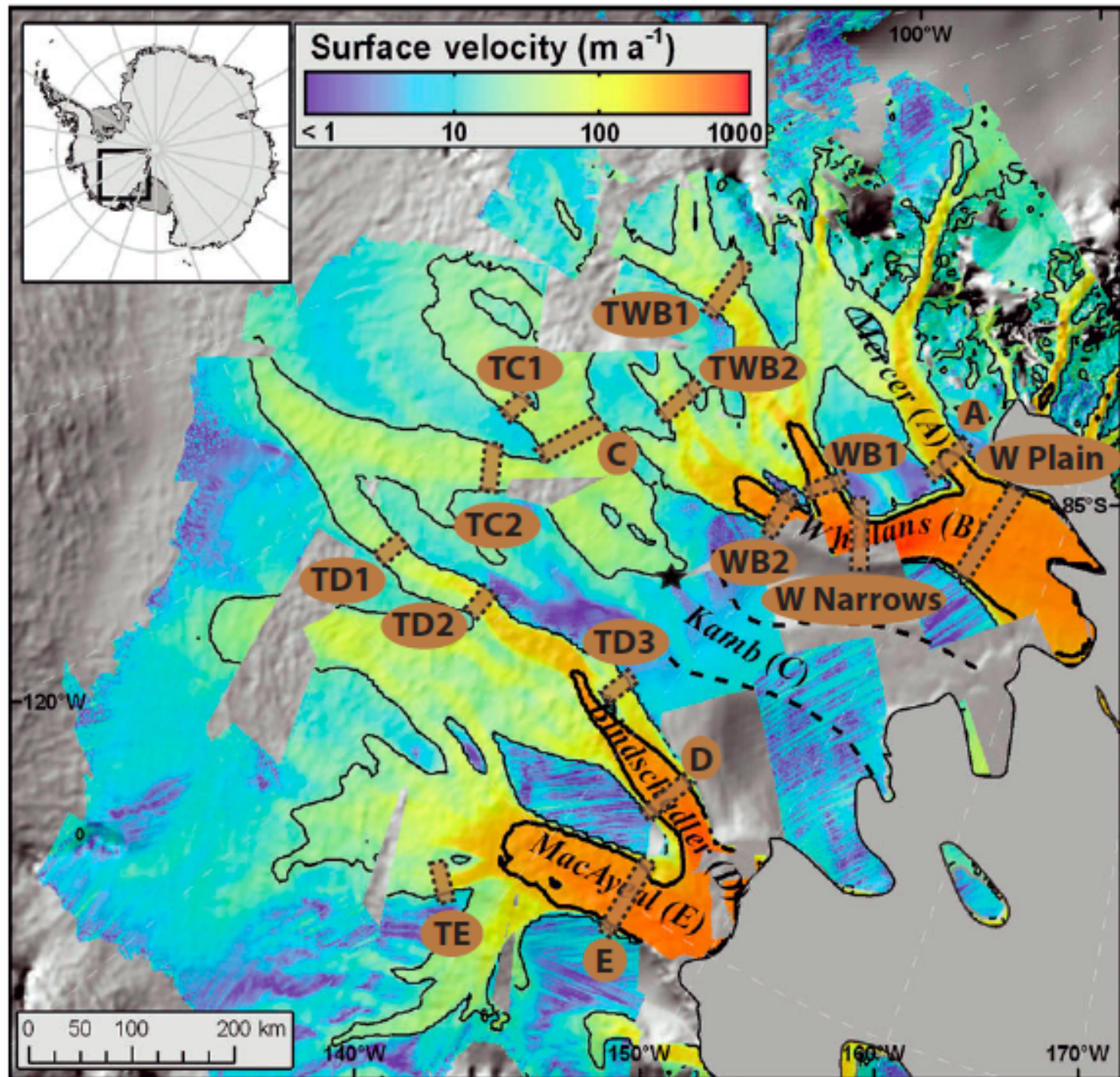
⇒ onset of melting ?

\*neglecting any variation in net axial force in sheet, roughly justified [Whillans and van der Veen, *J. Glac.*, 1993]

Our data set,  
to test  
concepts:

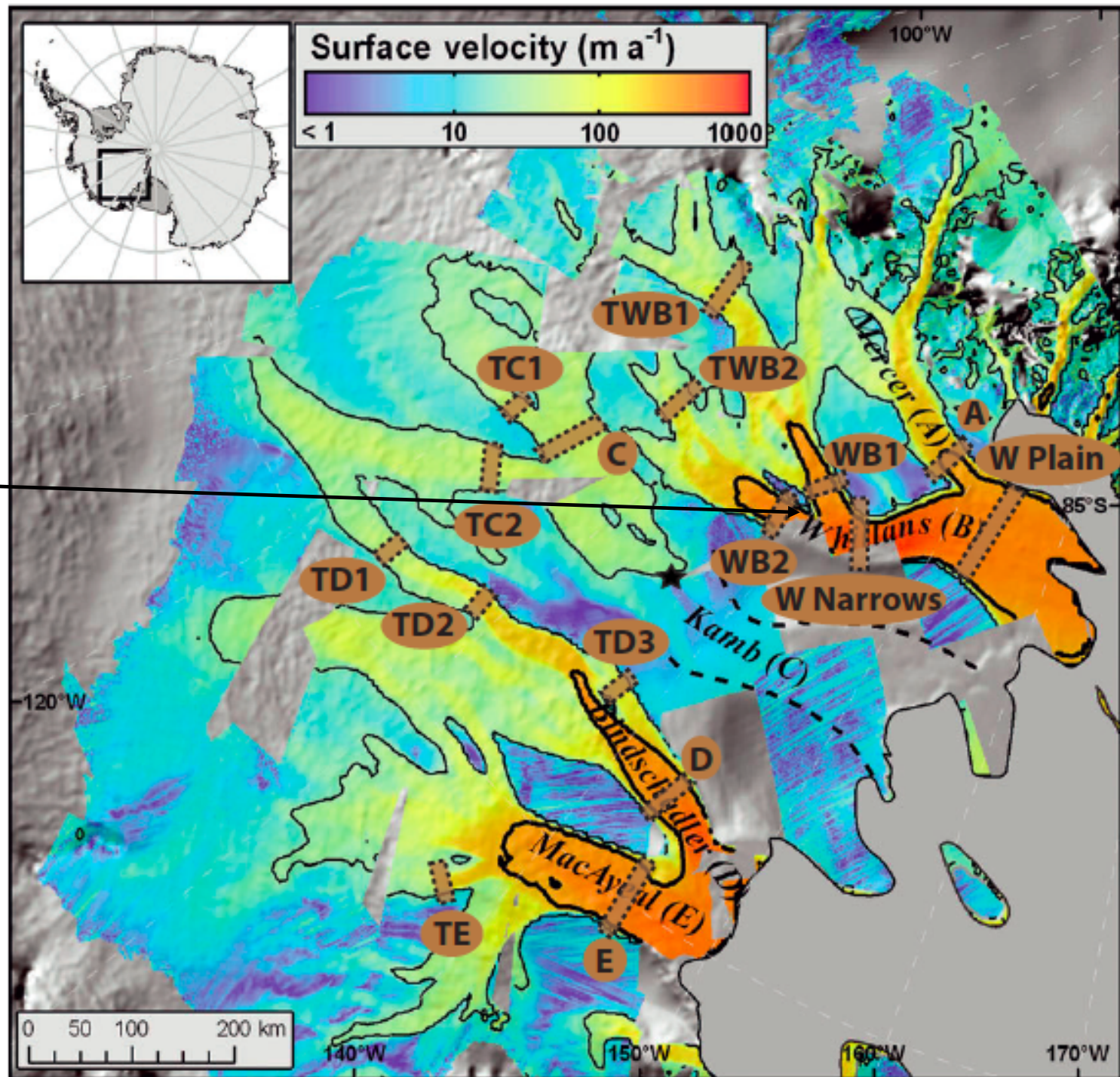
Sixteen  
ice stream  
traverses  
(dotted-lines)  
for velocity  
profiles.

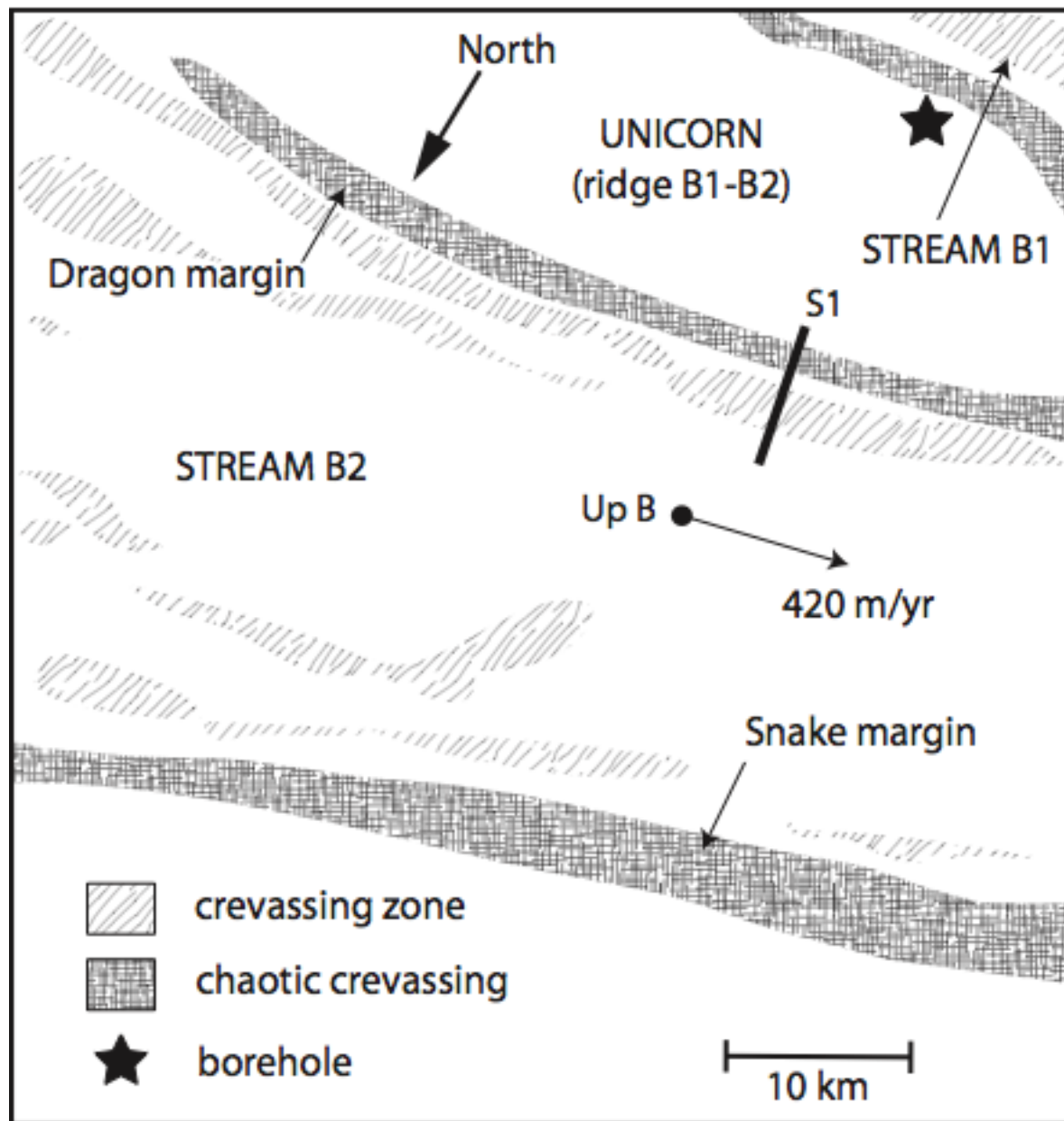
Characterized  
by Joughin et  
al. (*J. Geoph.  
Res.*, 2002)



Focus for  
most detailed  
study:

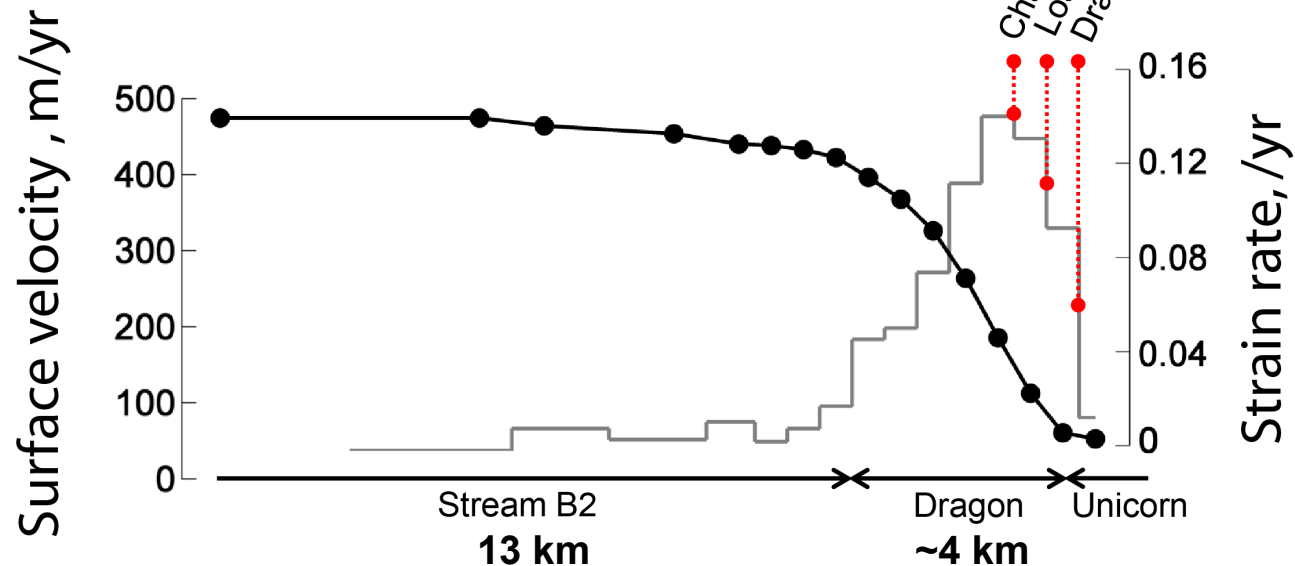
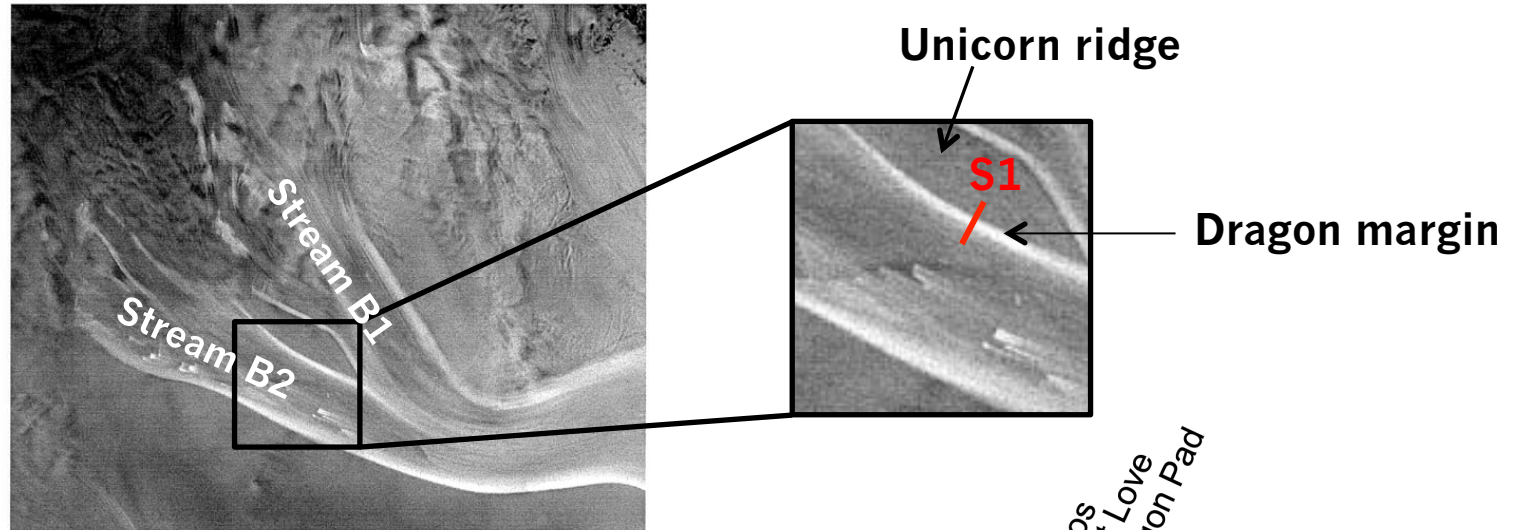
Dragon  
Margin —  
(very near to,  
but not the same as,  
the  
WB2 profile)





Echelmeyer and Harrison (1999)

# Dragon margin



**Table 1.** Parameters taken from *Joughin et al.* [2002] and used for margins of the profiles located in Figure 2. Profiles beginning with the letter T are made at the tributaries of the ice streams.

Ice Stream	Profile	$H$ (m)	$W$ (km)	$\dot{\gamma}_{lat}$ ( $10^{-2} \cdot \text{yr}^{-1}$ )
Mercer	A	1242	39	4.2
Whillans	WB1	1205	35	7.0
	WB2	985	34	9.5
	W Narrows	846	48	13.5
	W Plain	735	121	5.1
	TWB1	2188	25	3.8
	TWB2	1538	25	4.0
Kamb	C	1805	69	1.0
	TC1	1802	17	1.4
	TC2	2196	43	0.9
Bindschadler	D	888	55	5.8
	TD1	1952	24	2.5
	TD2	1412	35	5.4
	TD3	1126	21	2.2
MacAyeal	E	916	78	8.1
	TE	1177	19	5.5

**Lateral strain rate**

$\dot{\gamma}_{lat}$   
**here is an average  
over ~ 2 km width  
at Ice Stream margin**

# Thermo - mechanical properties of ice

Data fits as suggested by  
Cuffy & Paterson [2010]

Glen's flow law for ice  
(dislocation creep) :

$$\dot{\gamma} / 2 = A(T) \tau^3,$$

$$\tau = B(T) (\dot{\gamma} / 2)^{1/3}$$

$$\left( B(T) = [1 / A(T)]^{1/3} \right)$$

With diffusion creep too,

$$\dot{\gamma} / 2 = \tau / 2\eta(T, d_g) + A(T) \tau^3$$

Ice is still strong at  $T_{melt}$ :

For a given  $\dot{\gamma}$ :

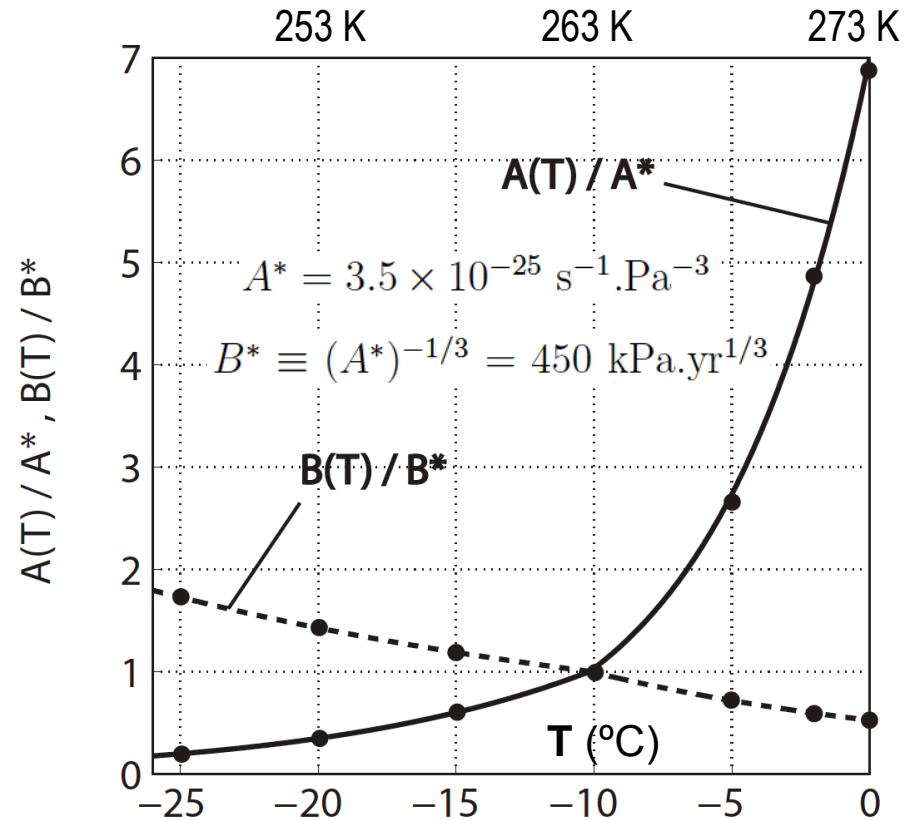
$$\tau_{T=0^\circ\text{C}} \approx 0.5 \times \tau_{T=-13^\circ\text{C}}$$

Thermal conductivity :

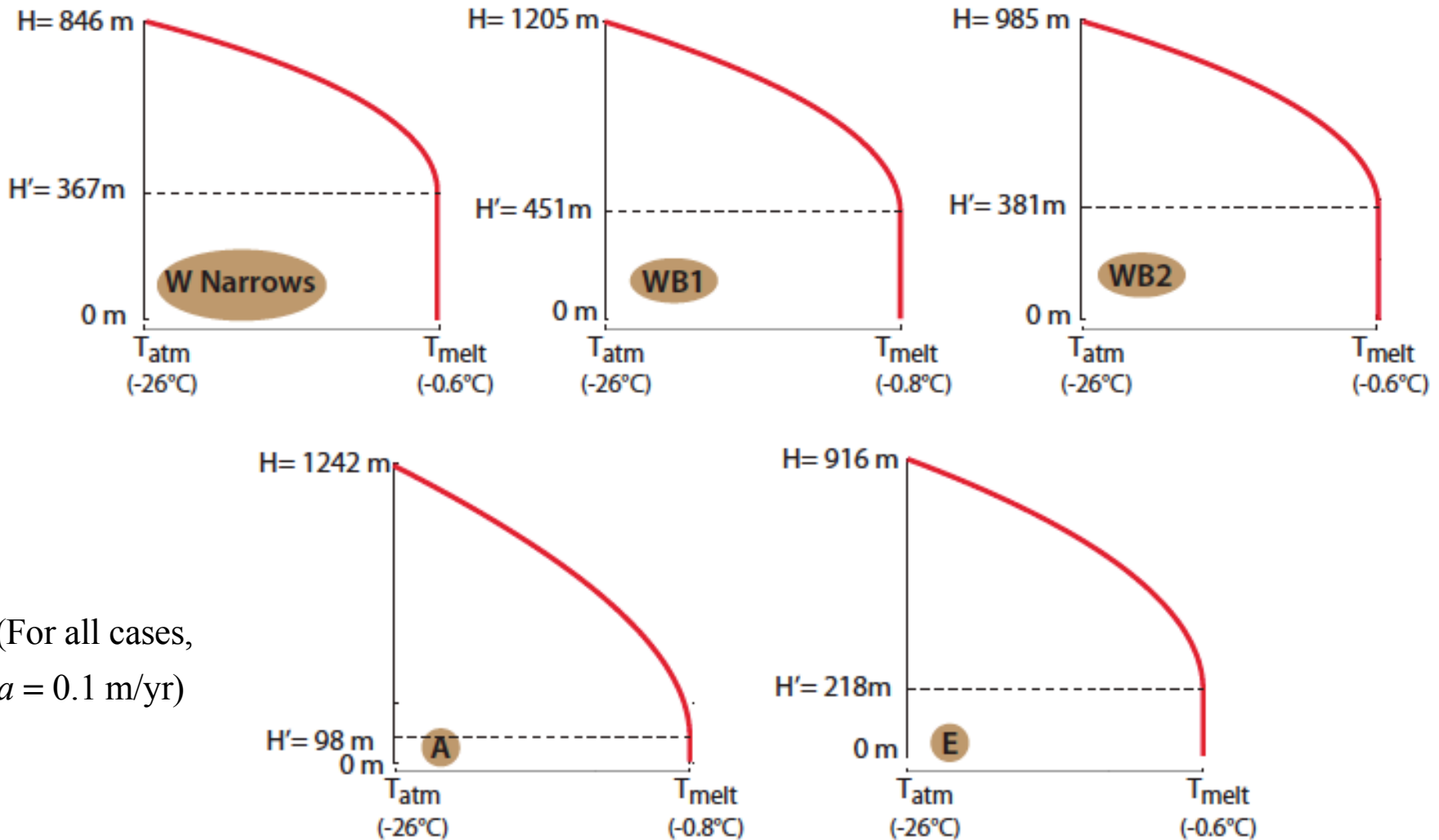
$$K(T) = 9.828 \frac{\text{J}}{\text{m s K}} \exp\left(-5.7 \times 10^{-3} \frac{T}{\text{K}}\right)$$

Specific heat :

$$C_i(T) = \left( 152.5 + 7.122 \frac{T}{\text{K}} \right) \frac{\text{J}}{\text{kg K}}$$



$$\frac{d}{dz} \left( K(T) \frac{dT}{dz} \right) + \left( \frac{1}{2A(T)} \right)^{1/3} (\dot{\gamma}_{lat})^{4/3} + \rho C_i(T) (az / H) \frac{dT}{dz} = 0 \quad \& \quad T \leq T_{melt} \Rightarrow$$



(For all cases,  
 $a = 0.1 \text{ m/yr}$ )

**Typical results, valid for nearly all stream margins:  
 Substantial thickness of temperate ice (i.e., at  $T = T_{melt}$ )  
 predicted at base of the ice sheet**

**2D, antiplane strain analysis :** velocity =  $u(y,z)\vec{e}_x$  ,  $\dot{\gamma} = \sqrt{\vec{\nabla}u \cdot \vec{\nabla}u}$

**Coupled non - linear Poisson equation system**

(for velocity  $u$  and temperature  $T$ ) :

$$\frac{\partial}{\partial y} \left( \frac{\tau(\dot{\gamma}, T)}{\dot{\gamma}} \frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left( \frac{\tau(\dot{\gamma}, T)}{\dot{\gamma}} \frac{\partial u}{\partial z} \right) = -\rho g S \quad \dot{\gamma} = \max \left( 2A(T)\tau^3, \tau / \eta(T) \right)$$

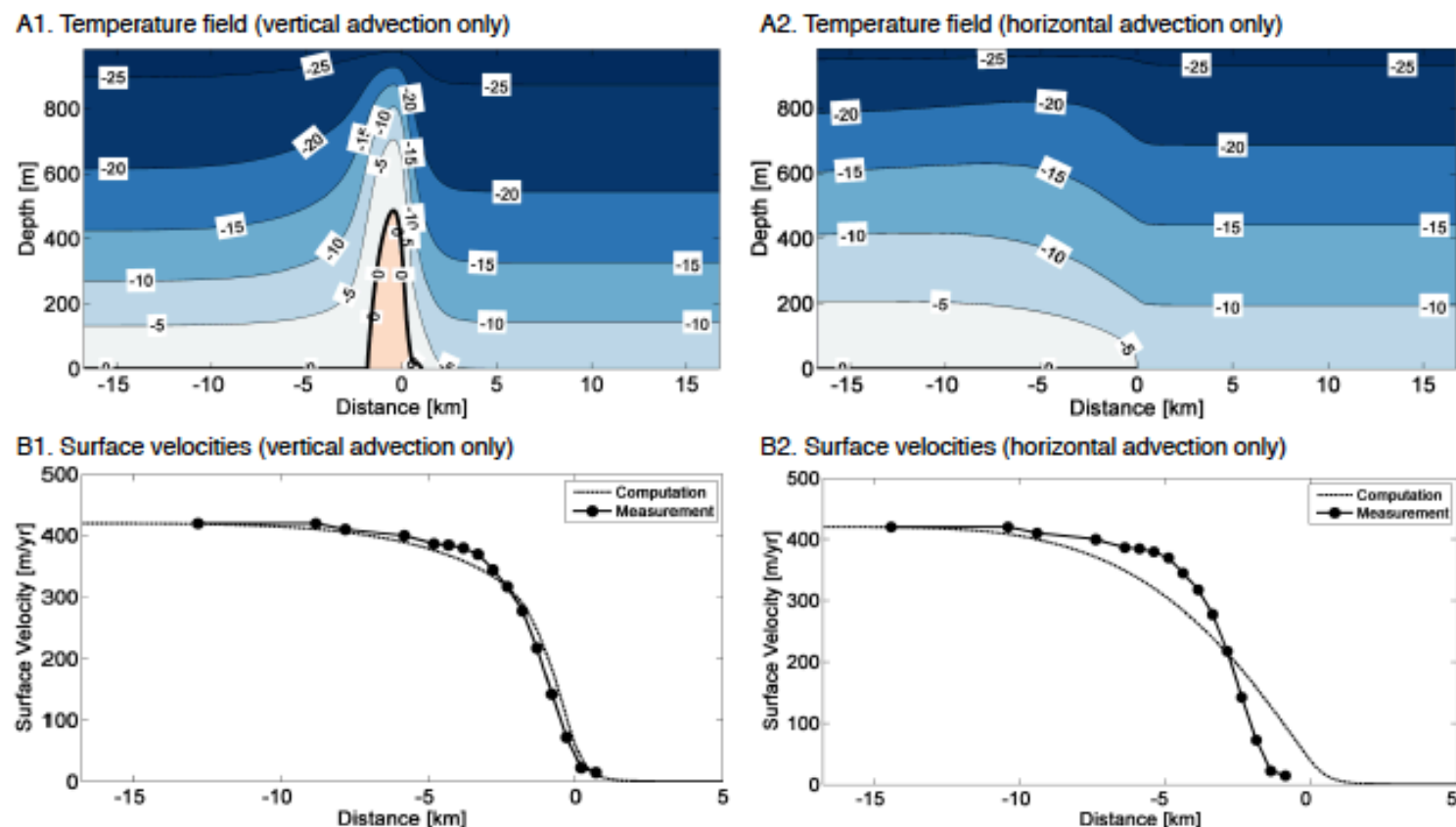
$$\begin{aligned} & \frac{\partial}{\partial y} \left( K(T) \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( K(T) \frac{\partial T}{\partial z} \right) \\ &= -[1 - \hat{H}(T - T_{melt})] \tau(\dot{\gamma}, T) \dot{\gamma} + \rho C(T) \left( v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} \right) \end{aligned}$$

(here,  $v$  &  $w$  are regarded as given, e.g., Zotikov's  $w = -az / H$ ,  $v = \text{const.}$ ,  $v_o$ )

**Computational Approach (Suckale, Platt, Perol, Rice [JGR, 2014]) :**

- **Multigrid methodology** for iterative solution of coupled nonlinear Poisson systems, embedding constraint  $T \leq T_{melt}$ .

(Suckale, Platt, Perol, and Rice, JGR 2014)



**Figure 8.** Temperature fields and surface velocities for Dragon margin when including only vertical advection (A1 and B1) with  $a = 0.1$  m/yr and only horizontal advection (A2 and B2) with  $v = -7.3$  m/yr, respectively. The best fitting basal stresses are  $\tau_{base} = 5.31$  kPa (A1 and B1) and  $\tau_{base} = 0.94$  kPa (A2 and B2), respectively. Both computations neglect surface crevassing.

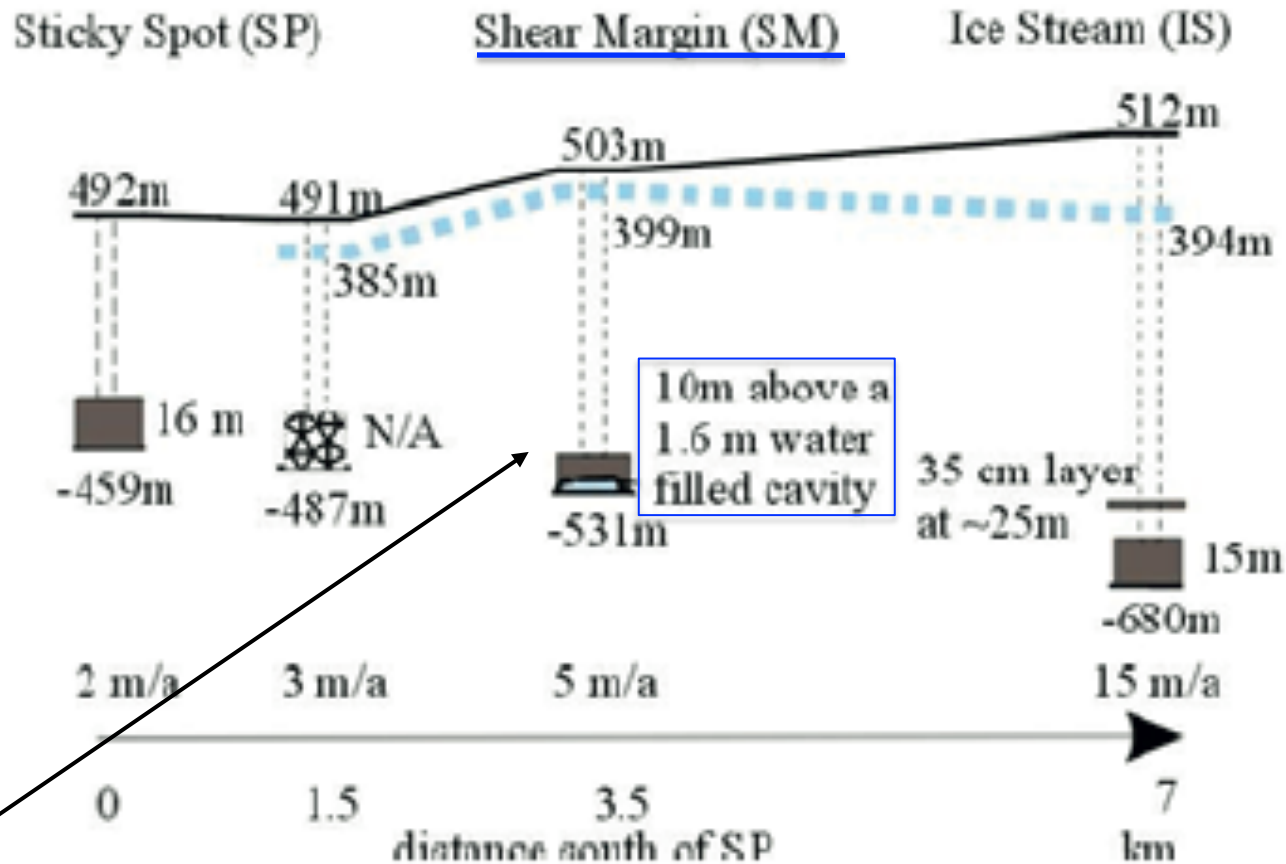
# Water drainage along glacier beds often creates channels

Natural example, shown in photo by Laura Kehrl: Rothlisberger Channel  
(but in a mountain glacier, with channel probably created by outburst underflooding):

Kennicott Glacier, Alaska



# Evidence of channel at margins

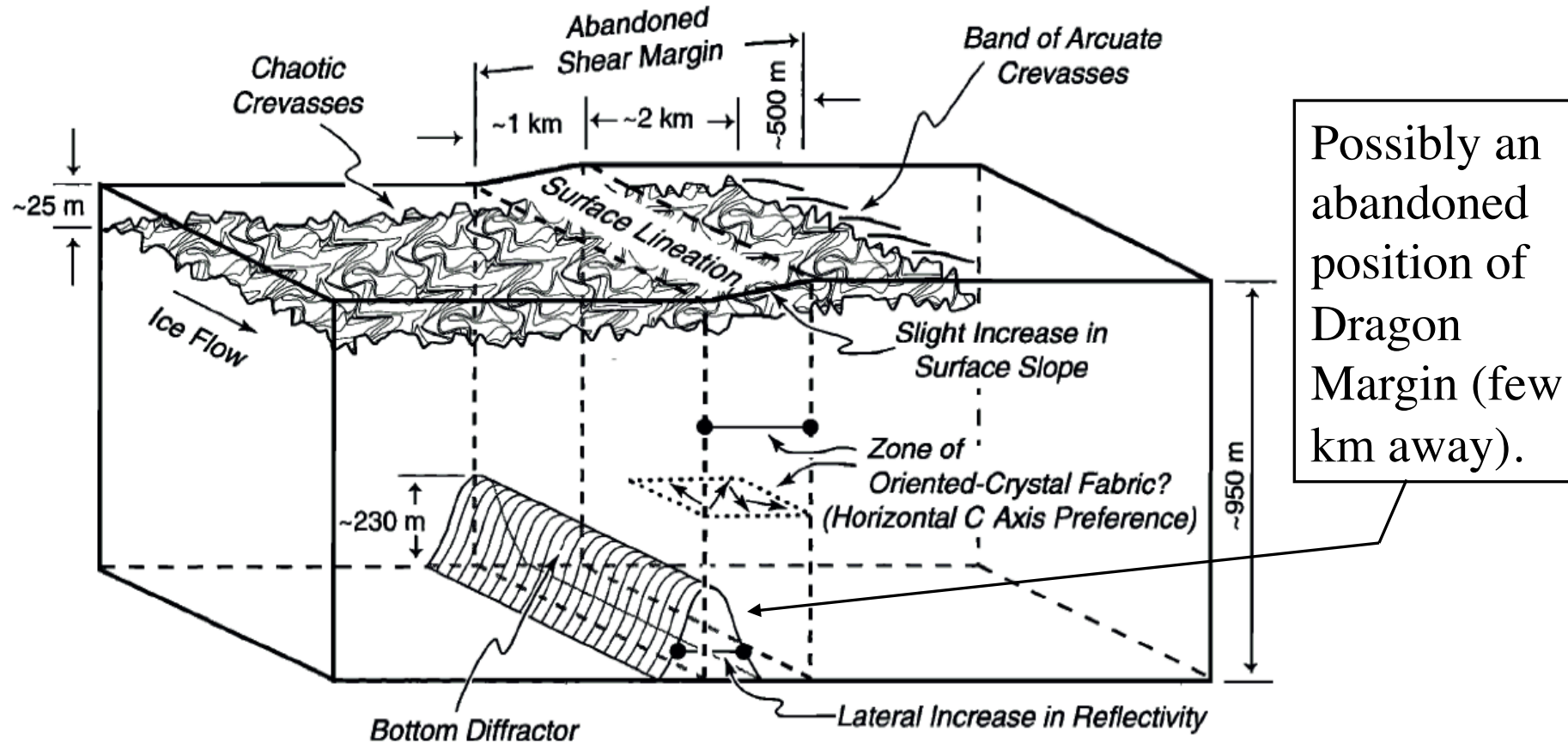


Borehole observation at the presently inactive shear margin of ***Kamb*** (C) ice stream:

- Found a 1.6 m tall water-filled cavity between 10 m of accreted ice and bed.
- Video of the borehole shows horizontal acceleration of particles sinking into the cavity, indicating flow of water within the cavity -- part of a channel?

[Modified from Vogel PhD (Thesis, 2004) and Vogel et al. (GRL, 2005).]

# Possible field evidence of internal melting at margins



- Clarke et al. [2000], in order to explain the bottom diffractors, have invoked **partial melting in temperate ice to a height of 230 m**, due to strain heating, among other possibilities (**entrained sediments, bottom crevasses**).
- Also, Clarke et al. noted a personal communication from H. Engelhardt (Caltech): Abnormal drill resistance encountered from  $\approx 56$  m above bed. Fresh scratches found on drill tip (assumed to be due to **entrained sediments**).

For  $Q_w = 100 \text{ km} \times 41 \text{ m}^3 / \text{m} \cdot \text{yr}$ , and  $S = 0.0012$ :

Clarke sub-glacial flooding range

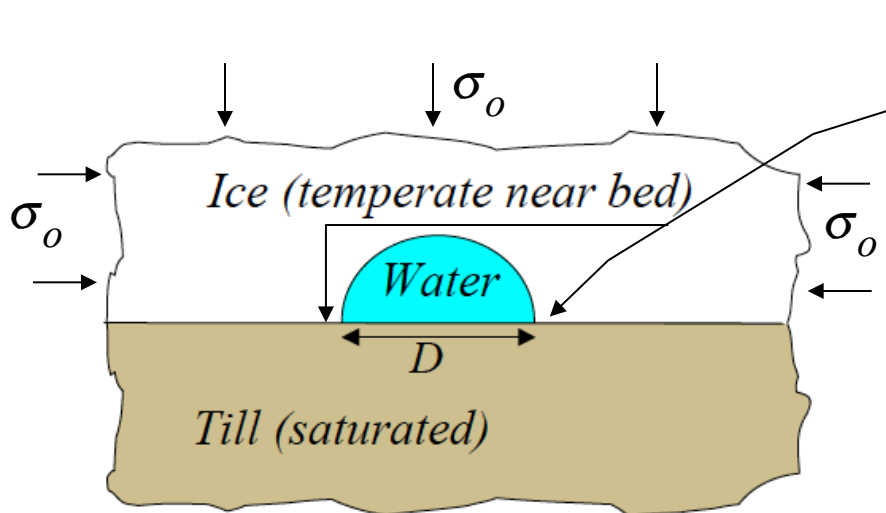
Assumed plausible here

Manning Coefficient, $n_M$ (s / m <sup>1/3</sup> )	0.01	0.02	0.03	0.04
Equivalent Nikuradse Roughness, $k$ (cm)	0.03	1.6	18.0	101.1
Channel Diameter, $D$ (m)	0.9	1.1	1.3	1.5
Effective Normal Stress at Channel Margin, $\sigma_{hoop} - p_{ch}$ (kPa)	369	310	280	261

# Rothlisberger-Shreve channel analysis

$$\sigma_{hoop} - p_{ch} = \frac{2}{3}(\sigma_o - p_{ch})$$

**Sensitivity :**  $Q_w \rightarrow 0.25Q_w \Rightarrow D \rightarrow 0.59D$ ,  $(\sigma_{hoop} - p_{ch}) \rightarrow 0.89(\sigma_{hoop} - p_{ch})$



Strength  $\tau_{ch} = f(\sigma_{hoop} - p_{ch})$

$$\approx 0.5(\sigma_{hoop} - p_{ch}) \approx 150 \text{ kPa}$$

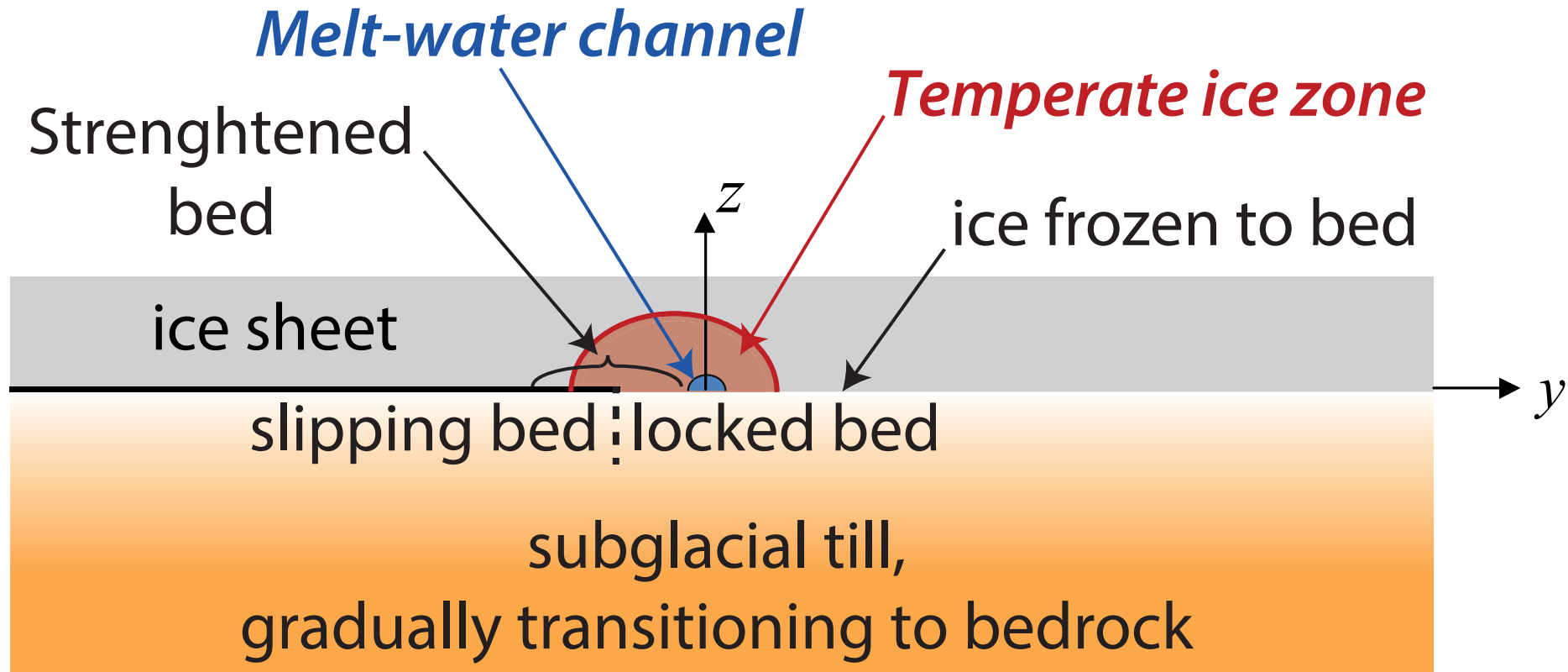
$$\therefore \tau_{ch} / \tau_{base} \approx 20 \text{ to } 45$$

For the 6 major streams,  $\tau_{ch} / \tau_{base}$  average = 32, and range = 12 to 56.

**So, a marginal drainage channel could be the source of enhanced basal resistance!**

(Perol, Rice, Platt and Suckale, AGU Dec. 2014)

*How subglacial hydrology can control  
the shear margin location of ice streams*



## Governing equations, mechanical-thermal-hydrologic model of ice stream

- **Mechanical model of an anti-plane shear flow driven by gravity**

$$\frac{\partial}{\partial y} \left( \frac{\tau(\dot{\gamma}, T)}{\dot{\gamma}} \frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left( \frac{\tau(\dot{\gamma}, T)}{\dot{\gamma}} \frac{\partial u}{\partial z} \right) + \rho_{ice} g S = 0, \quad \tau(\dot{\gamma}, T) = \min \left[ \left( \frac{\dot{\gamma}}{2A(T)} \right)^{1/3}, \eta(T) \dot{\gamma} \right]$$

- **Thermal model**

$$\frac{\partial}{\partial y} \left( K \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( K \frac{\partial T}{\partial z} \right) - \rho_{ice} C_{ice} \left( v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} \right) + [1 - H(T - T_{melt})] \tau \dot{\gamma} = 0$$

(we take  $v = 0$  (lateral advection neglected), and  $w = -az / H$  (Zotikov))

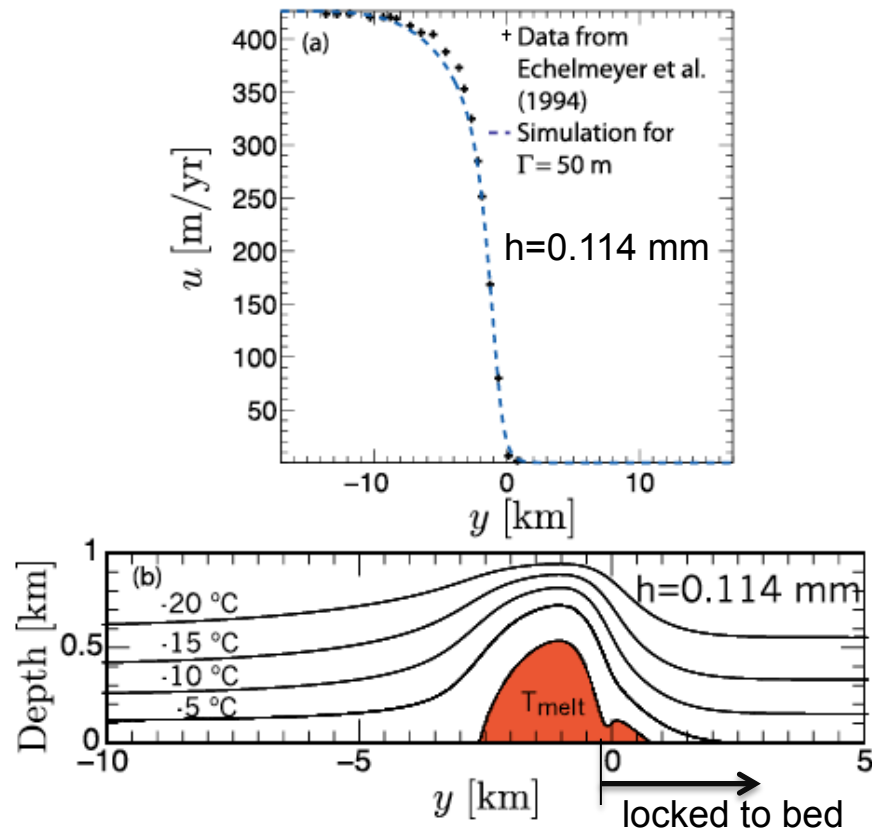
- **Subglacial hydrology model (Poiseuille flow in a thin water film)**

$$q_{melt} \equiv \frac{G_{geo} - G_{ice} + \tau_{base} u_b}{\rho_w L} = - \frac{\partial}{\partial y} \left( \frac{h^3}{12\mu_w} \frac{\partial p}{\partial y} \right) \quad (\text{we take thickness } h = \text{const.})$$

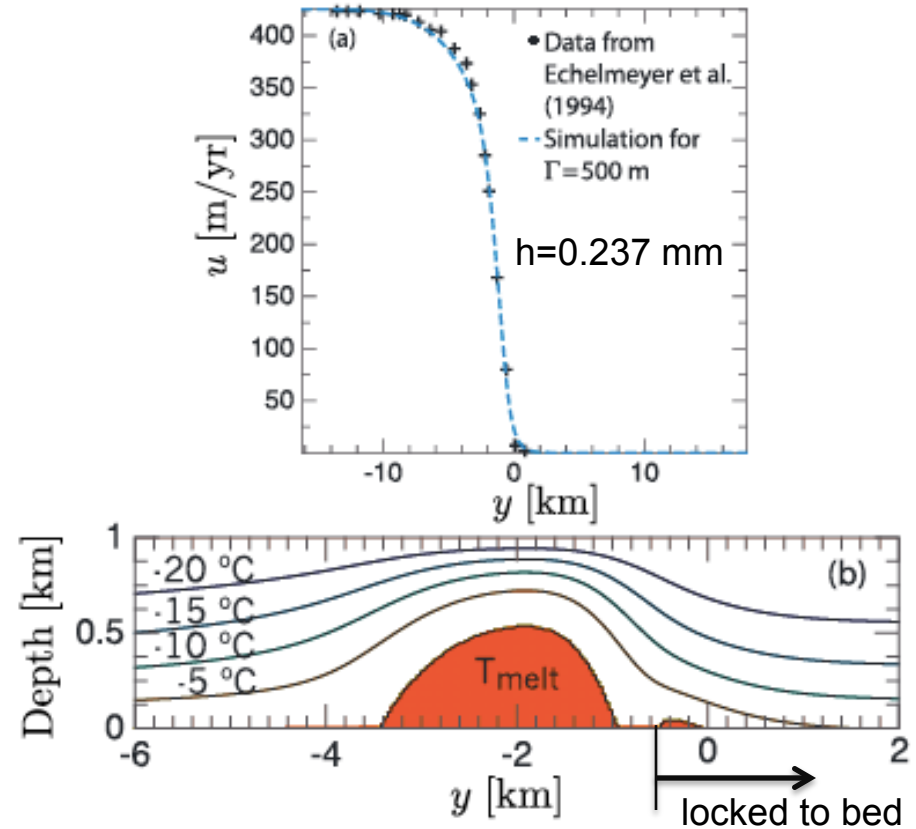
( $\tau_{base} = (\tau_{zx})_{z=0} = f \times (\rho_{ice} g H - p) + c$  = basal strength,  $u_b = u_{z=0}$  = basal sliding velocity)

**System solved using Finite Element procedure in COMSOL**

Ice surface deformation data is fit well by a range of models, with the slipping to locked transition occurring at 50 m to 500 m from the R-Channel (at  $y = 0$  here), corresponding to equivalent Poiseuille flow film thicknesses of 0.114 to 0.237 mm:



**Figure 5.** Results for a locking point 50 m away from the channel ( $\Gamma = 50$  m). (a) Numerical surface velocities plotted alongside the data from Echelmeyer et al. [1994]. (b) Temperature field. In red is the temperate ice zone in which the temperature is capped at the melting point. The basal shear stress is free of singularity for  $h = 0.114$  mm.



**Figure 7.** Results for a locking point 500 m away from the channel ( $\Gamma = 500$  m). (a) Numerical surface velocities compare with data from Echelmeyer et al. [1994]. (b) Temperature field. In red is the temperate ice zone in which the temperature is capped at the melting point. The basal shear stress is free of singularity for  $h = 0.237$  mm.